

Phase-space methods for quantum simulations

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The problem: complexity

How can we calculate quantum dynamics?

BEC: 10^{100000} states, 10^6 qubits in Hilbert space

- ✘ Real time path integrals **don't converge**
- ✘ Mean field methods **don't give quantum statistics**
- ✘ Direct computation needs **don't fit into memory!**

Phase-space representations

- ✓ Quantum dynamics \rightarrow stochastic motion + sampling
- ✗ **Classical** phase-space: **Wigner, P-, Q,** d dimensions
- ✓ **Quantum** phase-space: **Positive-P,** $2d$ dimensions
- ✓ Stochastic gauge: adds a weight to the trajectory

Phase-Space Representations

Expand the density matrix $\hat{\rho}$, using operators $\hat{\Lambda}(\vec{\lambda})$:

$$\hat{\rho} = \int P(\vec{\lambda}) \hat{\Lambda}(\vec{\lambda}) d\vec{\lambda}$$

Quantum dynamics \rightarrow Trajectories in $\vec{\lambda}$.

Different basis choice $\hat{\Lambda}(\vec{\lambda}) \rightarrow$ different representation

General M -mode Gaussian operator

Normally-ordered exponential of a general quadratic form in the $2M$ -vector mode operator $\delta\underline{\hat{a}} = (\hat{\mathbf{a}}, \hat{\mathbf{a}}^\dagger) - \underline{\alpha}$, where $\underline{\alpha}$ is a $2M$ -vector c-number and $\hat{\mathbf{a}}$ is the vector of annihilation operators. For algebraic reasons, it is useful to employ normal ordering, and to introduce a compact notation using a generalized covariance $\underline{\underline{\sigma}}$:

$$\hat{\Lambda}(\vec{\lambda}) = \frac{\Omega}{\sqrt{|\underline{\underline{\sigma}}|}} : \exp \left[-\delta\underline{\hat{a}}^\dagger \underline{\underline{\sigma}}^{-1} \delta\underline{\hat{a}}/2 \right] : .$$

In this case, the phase-space is described by the complex variables $\vec{\lambda} = (\Omega, \underline{\underline{\sigma}}) = (\Omega, \alpha)$.

What is the covariance?

$$\underline{\underline{\sigma}} = \begin{bmatrix} \mathbf{I} + \mathbf{n} & \mathbf{m} \\ \mathbf{m}^+ & \mathbf{I} + \mathbf{n}^T \end{bmatrix} \cdot$$

The representation phase space is $\vec{\lambda} = (\Omega, \alpha, \alpha^+, \mathbf{n}, \mathbf{m}, \mathbf{m}^+)$

- Ω = weight factor
- α, α^\dagger = amplitude
- \mathbf{n} = number correlation (complex $M \times M$ matrix).
- \mathbf{m}, \mathbf{m}^+ = squeezing (symmetric complex $M \times M$ matrices)

What are the moments physically?

$$\langle \hat{a}_i \rangle = \langle \Omega \alpha_i \rangle_P$$

$$\langle \hat{a}_i^\dagger \rangle = \langle \Omega \alpha_i^+ \rangle_P$$

$$\langle \hat{a}_i \hat{a}_j \rangle = \langle \Omega (\alpha_i \alpha_j + m_{ij}) \rangle_P$$

$$\langle : \hat{a}_i \hat{a}_j^\dagger : \rangle = \langle \Omega (\alpha_i \alpha_j^+ + n_{ij}) \rangle_P$$

$$\langle \hat{a}_i^\dagger \hat{a}_j^\dagger \rangle = \langle \Omega (\alpha_i^+ \alpha_j^+ + m_{ij}^+) \rangle_P \cdot$$

Standard Phase-space Representations

What about the text-book phase-space representations?

<i>Property:</i> Repn.	<i>Variance</i> (n)	<i>Operator</i> <i>Order</i>	Phase-space
P	0	Normal	Classical
W	1/2	Symmetric	Classical
Q	1	Antinormal	Classical
+P	0	Normal	Classical $\times 2$
G	n	Any	(Classical) ²

TYPES OF PROBLEM

There are three main types of problems studied:

- ✓ Master equations - damping + coherent evolution
- ✓ Canonical ensembles - 'imaginary time' thermal equilibrium
- ✓ Quantum dynamics - purely coherent nonlinear evolution

OUTLINE

1. **PROBLEM:** $\partial \hat{\rho} / \partial t = \hat{L}[\hat{\rho}]$
2. Define: p -dimensional complex space: $\vec{\lambda} = (\Omega, \alpha)$
3. **Basis** $\hat{\Lambda}(\vec{\lambda})$: $\hat{\rho} = \int P(\vec{\lambda}) \hat{\Lambda}(\vec{\lambda}) d^{2p} \vec{\lambda}$
4. Identities: $\partial \hat{\rho} / \partial t = \int P(\vec{\lambda}) \mathcal{L}_A \hat{\Lambda}(\vec{\lambda}) d^{2p} \vec{\lambda}$
5. **Diffusion and drift gauge:** $\mathbf{g}^d(\alpha), \mathbf{g}(\alpha)$
6. **STOCHASTICS:**
 $d\Omega / \partial t = \Omega [U + \mathbf{g} \cdot \zeta]$
 $d\alpha / \partial t = \mathbf{A} + \mathbf{B}(\zeta - \mathbf{g})$

CENTRAL RESULT

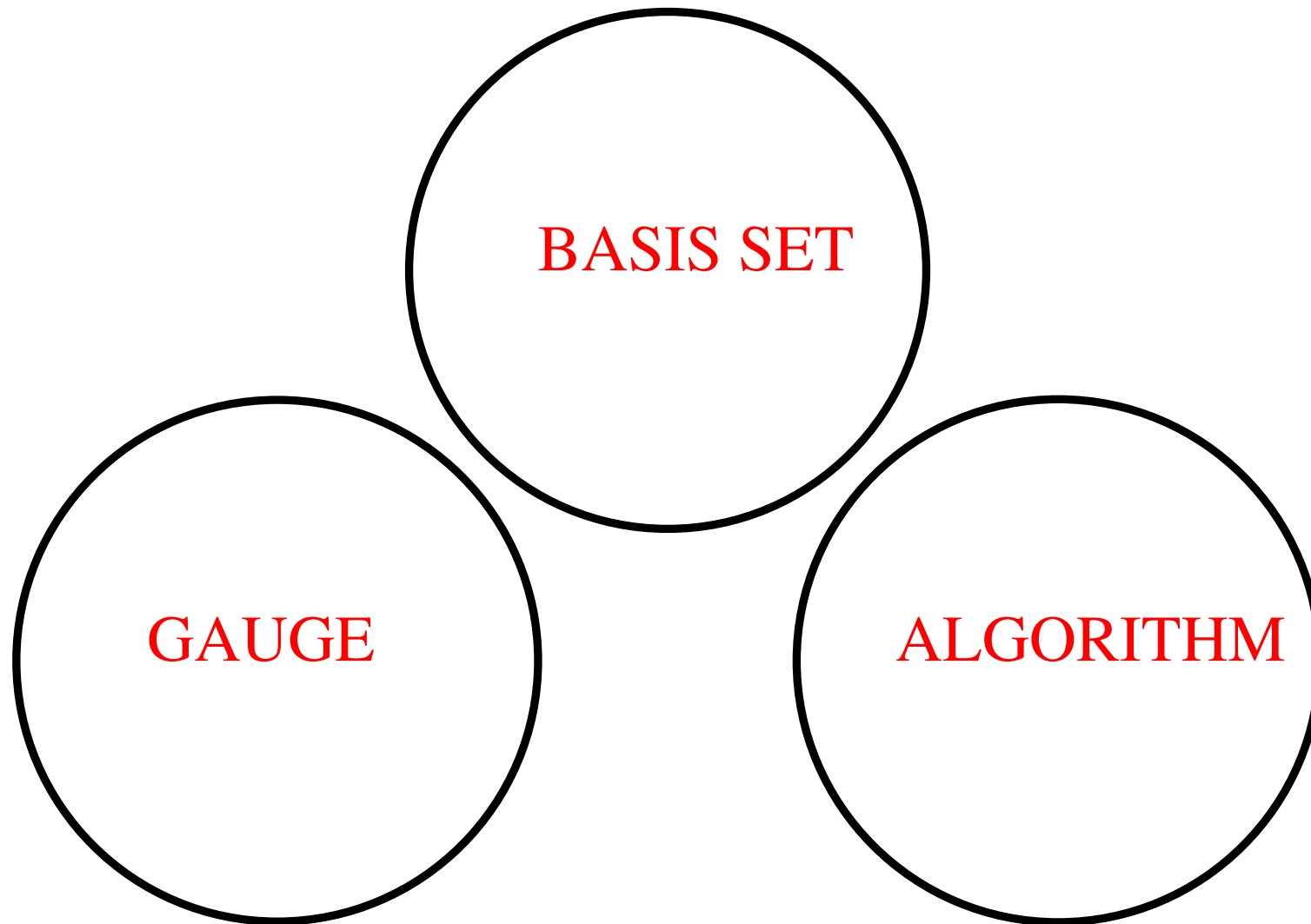
Stochastic equations: **trajectory** α , quantum **amplitude** Ω :

Gauge $d\Omega/\partial t = \Omega [U dt + \mathbf{g}(\alpha) \cdot \zeta(t)]$

Trajectory $d\alpha/\partial t = \mathbf{A} + \mathbf{B}[\zeta(t) - \mathbf{g}(\alpha)]$

- Noise correlations: $\langle \zeta_i(t) \zeta_j(t') \rangle = \delta_{ij} \delta(t - t')$
- Gauges chosen freely to optimize simulations
- Provided no boundary terms, all gauges EQUIVALENT
- Works for either bosons or fermions
- Sign changes for fermions and NO amplitude terms

COMPUTATIONAL STRATEGIES



League Table of Representations

How do the known phase-space representations compare?

<i>Property:</i> Repn.	<i>Finite?</i>	<i>2nd</i> <i>Order</i>	<i>Pos.</i> <i>Def.?</i>	<i>SDE ?</i>	<i>Stable?</i>
P	No	-	-	-	-
W	Yes	No	-	-	-
R	Yes	Yes	No	-	-
Q	Yes	Yes	Yes	No	-
+P	Yes	Yes	Yes	Yes	No
G	Yes	Yes	Yes	Yes	Yes

I: CANONICAL BOSE-HUBBARD MODEL

Nonlinear interactions at each site + linear interactions coupling different sites:

- $\hat{H}(\mathbf{a}, \mathbf{a}^\dagger) = \hbar \left[\sum \sum \omega_{ij} a_i^\dagger a_j + \sum : \hat{n}_j^2 : \right]$.
- ω_{ij} - nonlocal coupling, includes chemical potential.
- Boson number: $\hat{n}_i = a_i^\dagger a_i$.
- Grand canonical ensemble: $\hat{\rho}_u = e^{-(\hat{H} - \mu \hat{N})/k_B T} = e^{-\hat{K} \tau}$.

Stochastic gauge equations

Choose $n = 0$ then \rightarrow Imaginary time Gross- Pitaevskii equations with **weighting** and **quantum noise**:

$$\frac{d\alpha_i}{d\tau} = - \left[\alpha_i^\dagger \alpha_i + i g_i \right] \alpha_i - \sum_{j=1}^M \omega_{ij} \alpha_j / 2 + i \alpha_i \zeta_i(\tau)$$

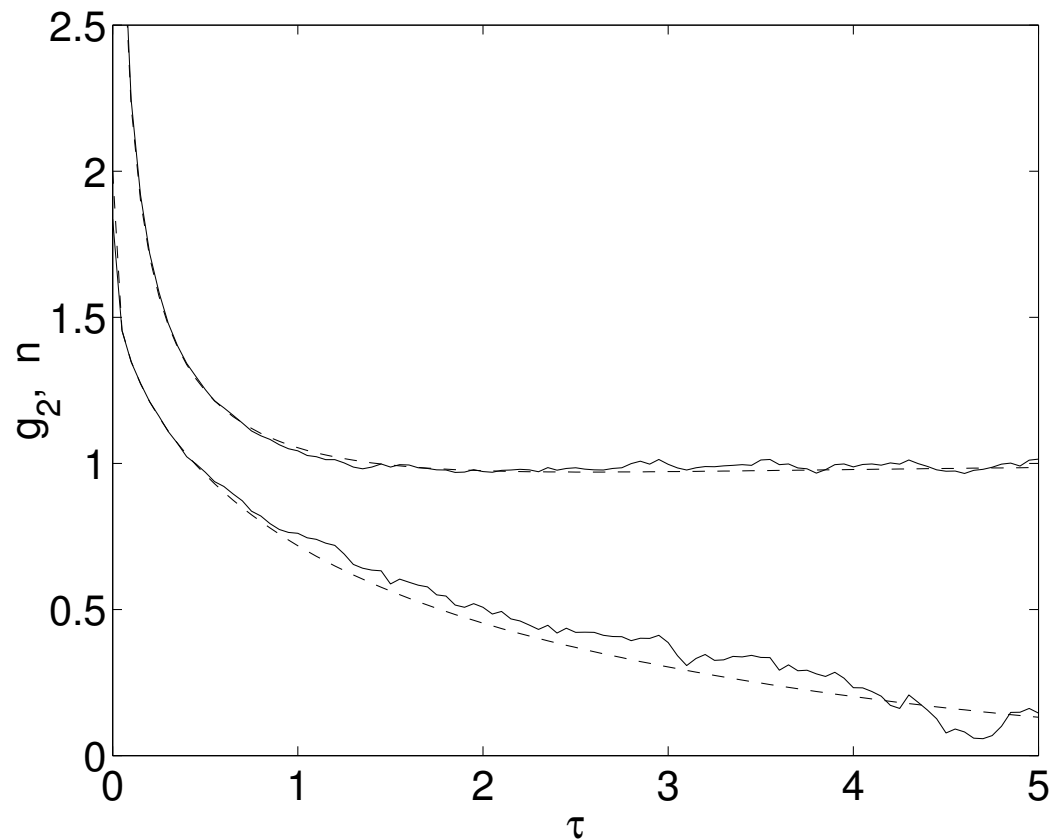
$$\frac{d\Omega}{d\tau} = \left[-K(\tau) d\tau + \sum_{i=1}^M g_i \zeta_i(\tau) \right] \Omega$$

STABILISING GAUGE: $g_j = i(\text{Re}(\alpha_j^\dagger \alpha_j) - |\alpha_j^\dagger \alpha_j|)$,

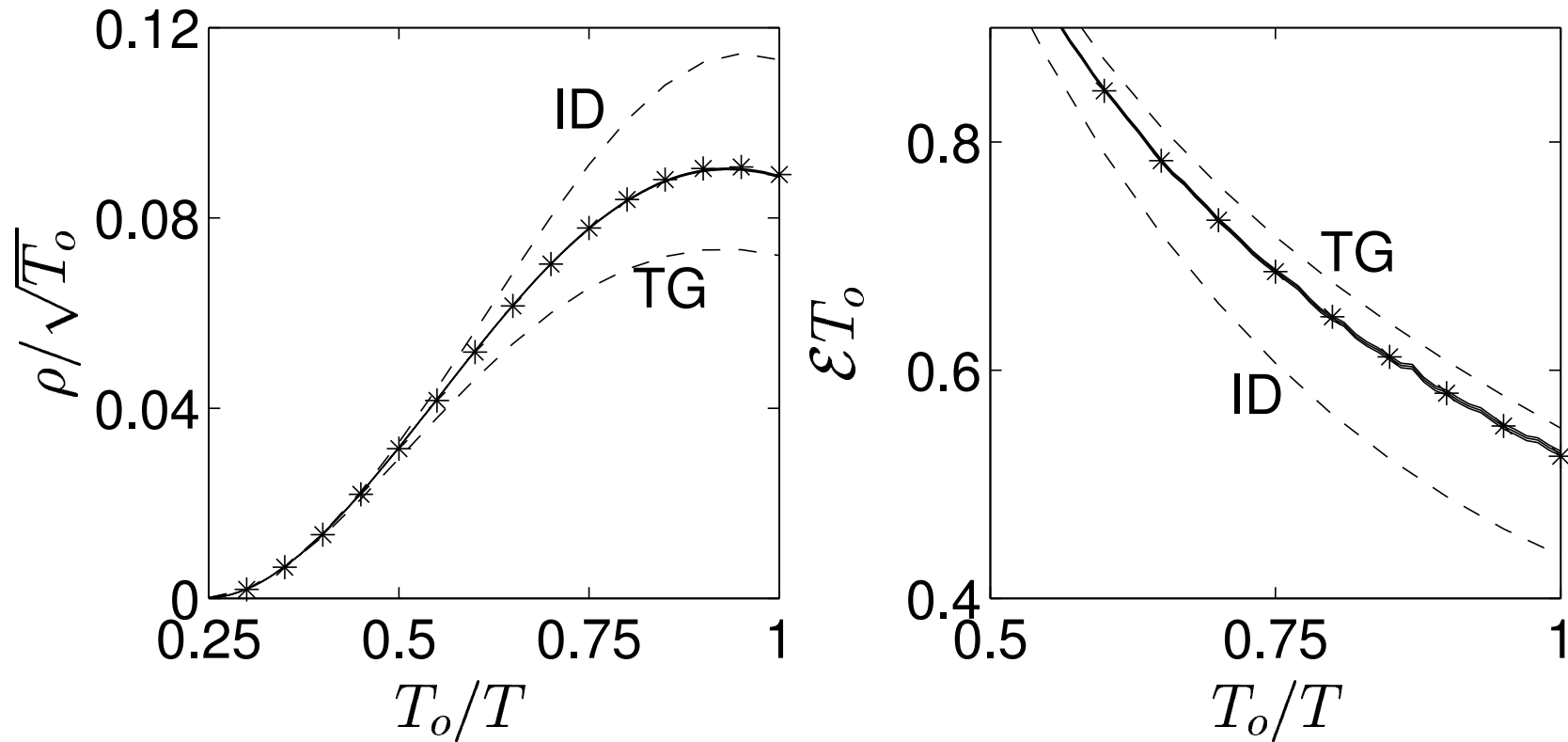
Single-well, interacting case

Antibunching: single Bose mode,

NOTE: τ = Inverse temperature



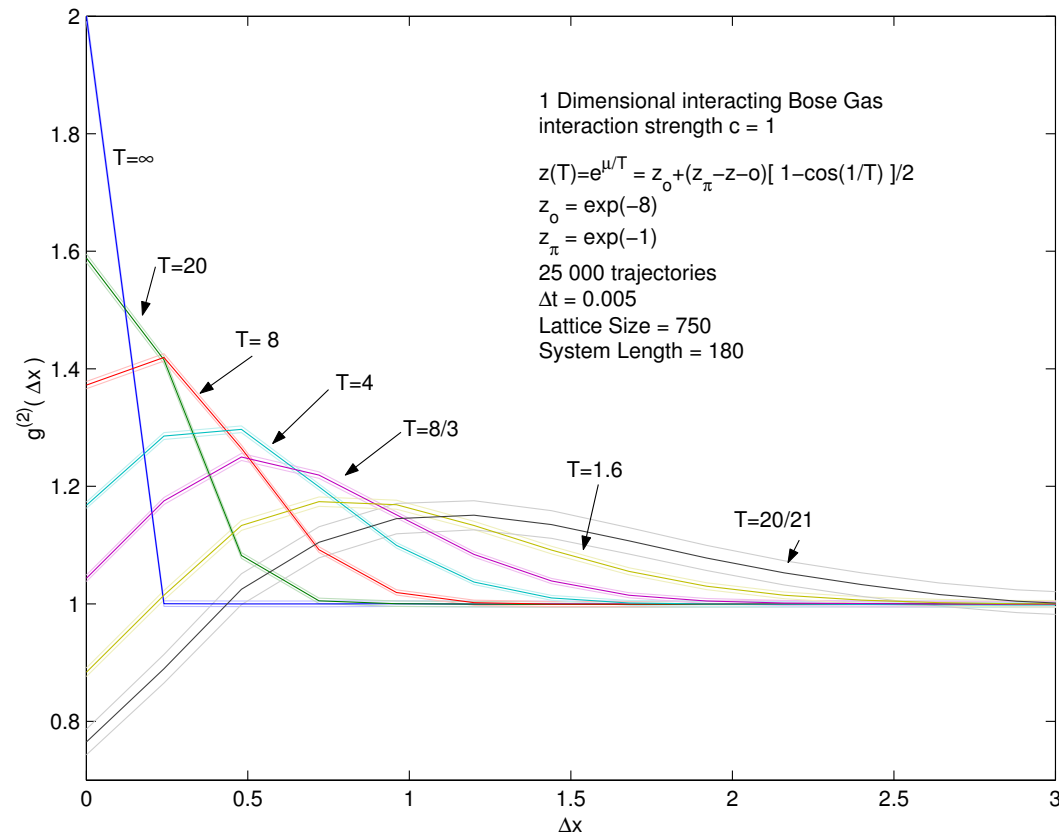
Simulations vs exact energy and density



- ✓ Complete agreement with exact solutions
 - at all temperatures calculated!

Spatial correlations

Spatial correlations, $g^{(2)}(x)$ can be calculated from gauge simulations:



II: REAL-TIME BOSE-HUBBARD MODEL

- $n = 0$ (*positive-P*) \rightarrow *Real time* Gross- Pitaevskii equations with **quantum noise**:

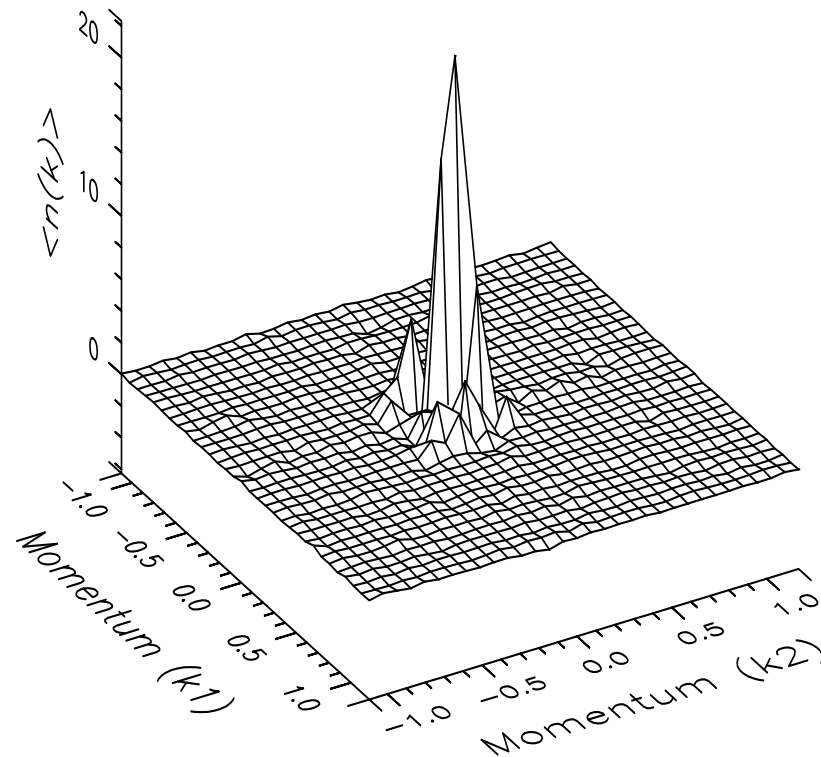
$$\frac{d\alpha_i}{dt} = -i \left[n_i + \sum_{j=1}^{2M} \omega_{ij} \alpha_j + \sqrt{i} \zeta_i(\tau) \right] \alpha_i$$

- $n = 1/2$ (*Wigner*) \rightarrow *Approximate* Gross- Pitaevskii equations with $\langle \alpha_i^{0*} \alpha_i^0 \rangle = 1/2$

$$\frac{d\alpha_i}{dt} = -i \left[n_i + \sum_{j=1}^{2M} \omega_{ij} \alpha_j \right] \alpha_i$$

BEC evaporative cooling (1998): 10^5 qubits

$t = 100$



III: NONCLASSICAL QUANTUM DYNAMICS:

Coherent molecular down-conversion

- Coherent process of molecular dissociation
- Overall effective Hamiltonian term in one dimension of

$$\hat{H} = \hat{H}_0 - i \frac{\hbar \chi(t)}{2} \int dx \left[e^{i\omega t} \hat{\Psi}_2^\dagger \hat{\Psi}_1^2 - e^{-i\omega t} \hat{\Psi}_2 \hat{\Psi}_1^{\dagger 2} \right],$$

- $\chi(t)$ is the bare atom-molecule coupling; ω is a detuning

Stochastic Equations

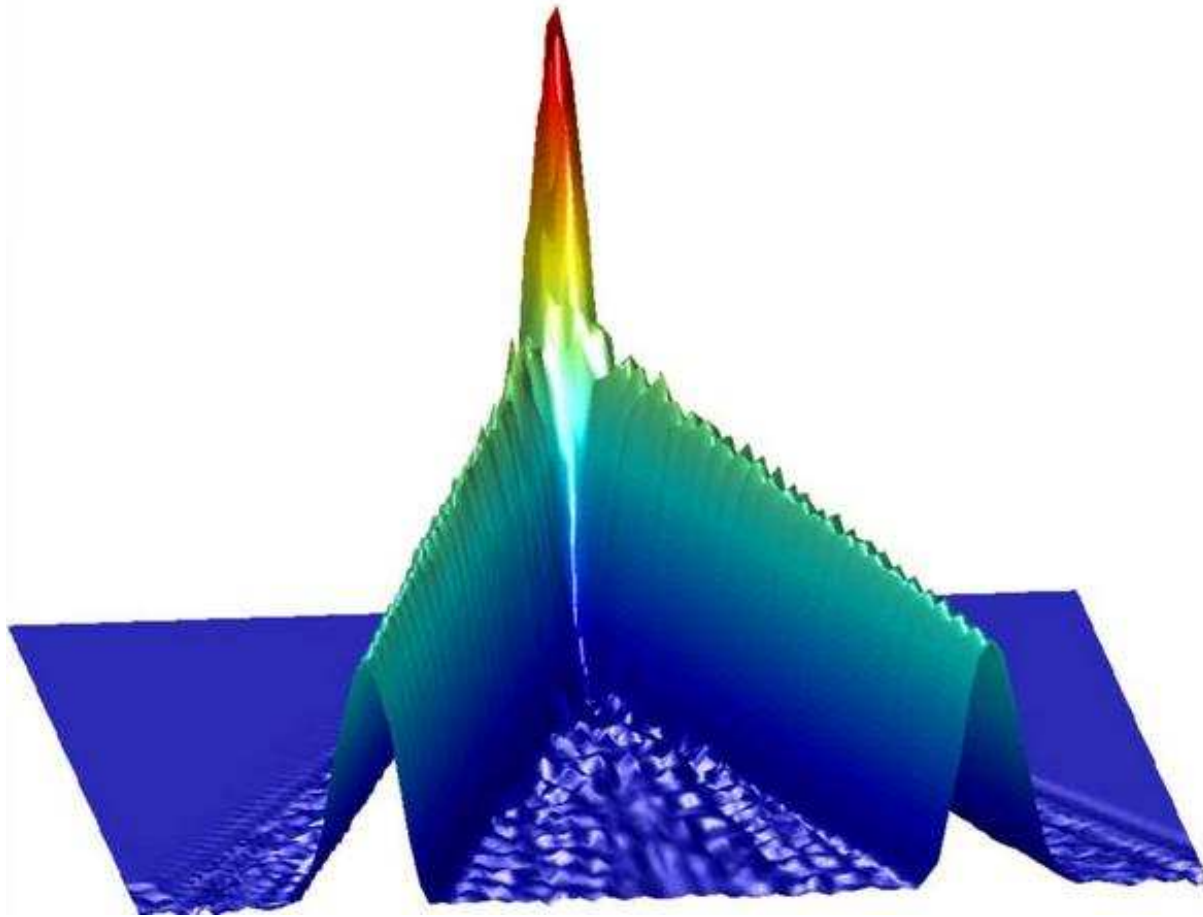
$$\frac{\partial \psi_1}{\partial \tau} = i \frac{\partial^2 \psi_1}{\partial \xi^2} - (\gamma + i\delta) \psi_1 + \kappa \psi_2 \psi_1^+ + \sqrt{\kappa \psi_2} \eta_1,$$

$$\frac{\partial \psi_1^+}{\partial \tau} = -i \frac{\partial^2 \psi_1^+}{\partial \xi^2} - (\gamma - i\delta) \psi_1^+ + \kappa \psi_2^+ \psi + \sqrt{\kappa \psi_2} \eta_1^+,$$

$$\frac{\partial \psi_2}{\partial \tau} = \frac{i}{2} \frac{\partial^2 \psi_2}{\partial \xi^2} - i\nu(\xi, \tau) \psi_2 - \frac{\kappa}{2} \psi_1^2 + \sqrt{-i\nu \psi_2} \eta_2,$$

$$\frac{\partial \psi_2^+}{\partial \tau} = -\frac{i}{2} \frac{\partial^2 \psi_2^+}{\partial \xi^2} + i\nu(\xi, \tau) \psi_2^+ - \frac{\kappa}{2} \psi_1^{+2} + \sqrt{i\nu \psi_2^+} \eta_2^+.$$

Twin atom correlations



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