

Disordered lattice gases

+ correlations in BEC

J.J. Arlt

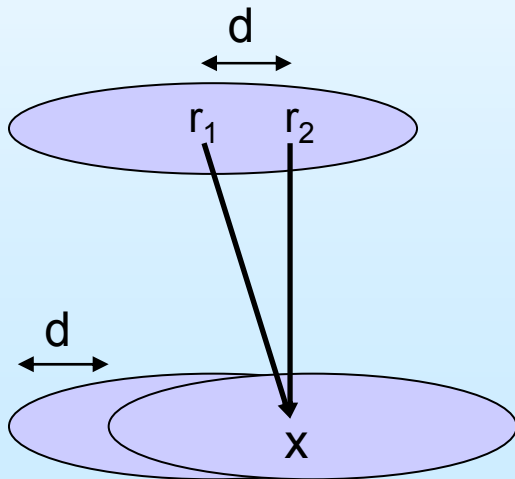
T. Schulte, S. Drenkelforth, W. Ertmer
K. Sacha, J. Zakrzewski und M. Lewenstein

1st order correlation function:

$$g^{(1)}(\mathbf{r}_1, \mathbf{r}_2) = \frac{\langle \hat{\psi}^\dagger(\mathbf{r}_1) \hat{\psi}(\mathbf{r}_2) \rangle}{\sqrt{\langle \hat{\psi}^\dagger(\mathbf{r}_1) \hat{\psi}(\mathbf{r}_1) \rangle \langle \hat{\psi}^\dagger(\mathbf{r}_2) \hat{\psi}(\mathbf{r}_2) \rangle}}$$

→ $\langle I(x) \rangle$

Contrast:

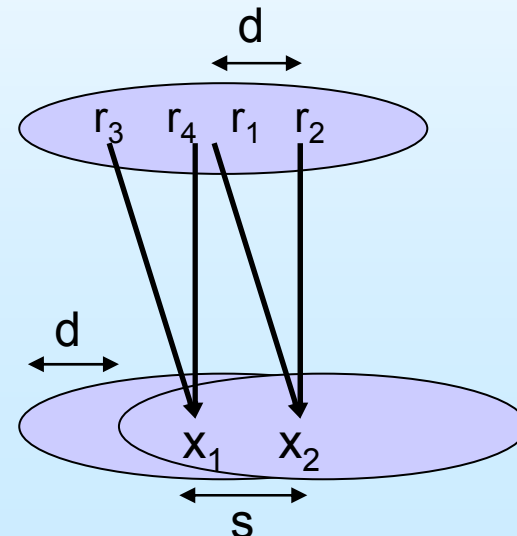


2nd order correlation function:

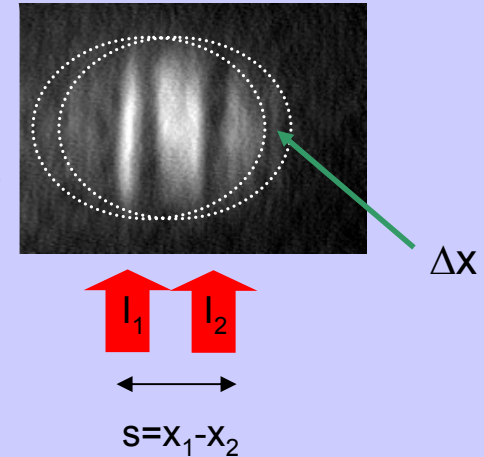
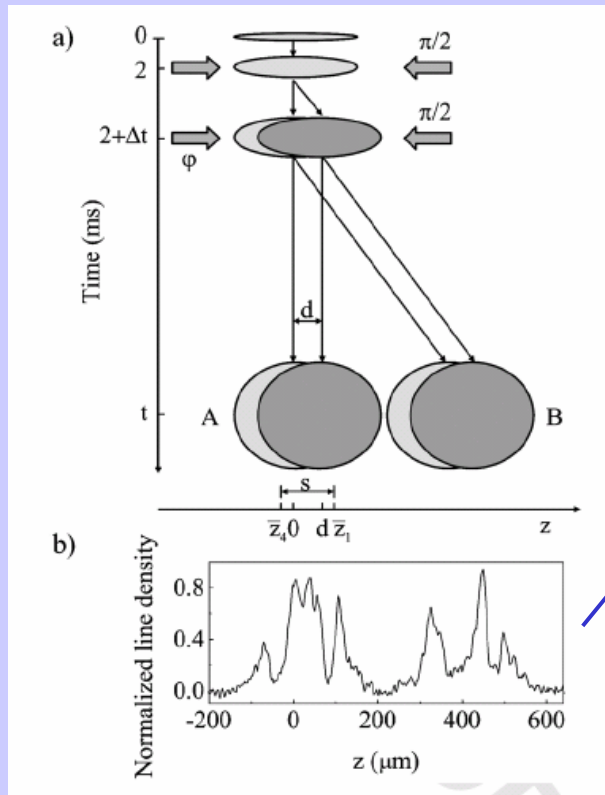
$$g^{(2)}(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \mathbf{r}_4) = \frac{\langle \hat{\psi}^\dagger(\mathbf{r}_1) \hat{\psi}^\dagger(\mathbf{r}_2) \hat{\psi}(\mathbf{r}_3) \hat{\psi}(\mathbf{r}_4) \rangle}{\sqrt{\prod_{i=1}^4 \langle \hat{\psi}^\dagger(\mathbf{r}_i) \hat{\psi}(\mathbf{r}_i) \rangle}}$$

→ $\langle I(x_1) I(x_2) \rangle$

Intensity correlations:



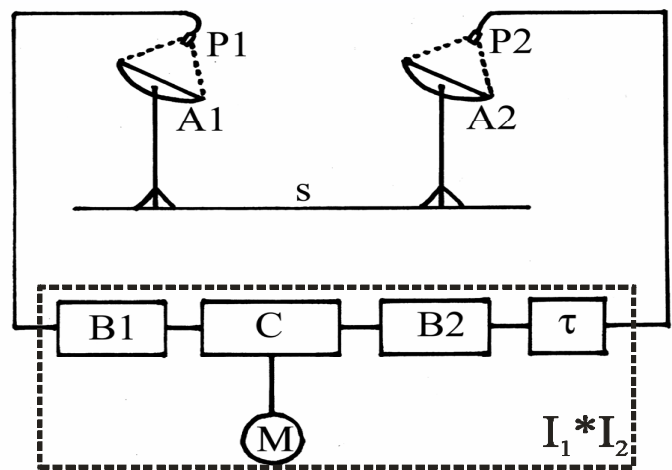
Interferometric scheme



$$\gamma_f^{(2)}(x_1, x_2, \Delta x) = \frac{\langle (\hat{I}_1 - \langle \hat{I}_1 \rangle) (\hat{I}_2 - \langle \hat{I}_2 \rangle) \rangle}{\sqrt{\langle (\hat{I}_1 - \langle \hat{I}_1 \rangle)^2 \rangle \langle (\hat{I}_2 - \langle \hat{I}_2 \rangle)^2 \rangle}}$$

The normalised intensity correlation function can be measured !

Hanbury-Brown Twiss



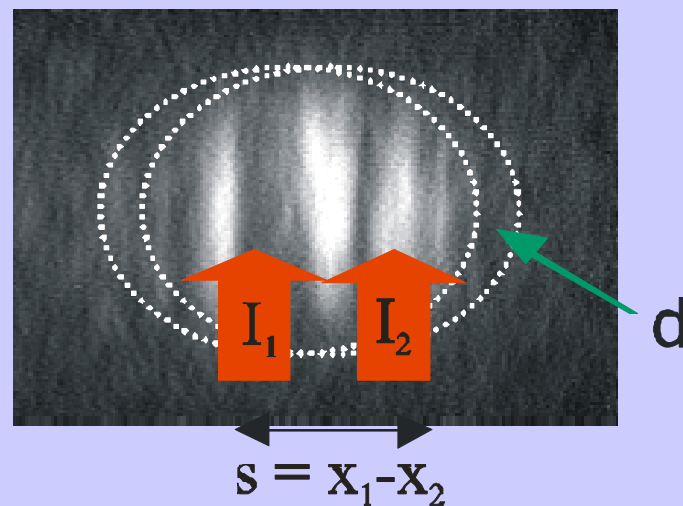
Measurement of intensity correlations

➡ transverse coherence length

➡ star diameters

Advantage:
insensitive to atmospheric fluctuations

Interferometric scheme



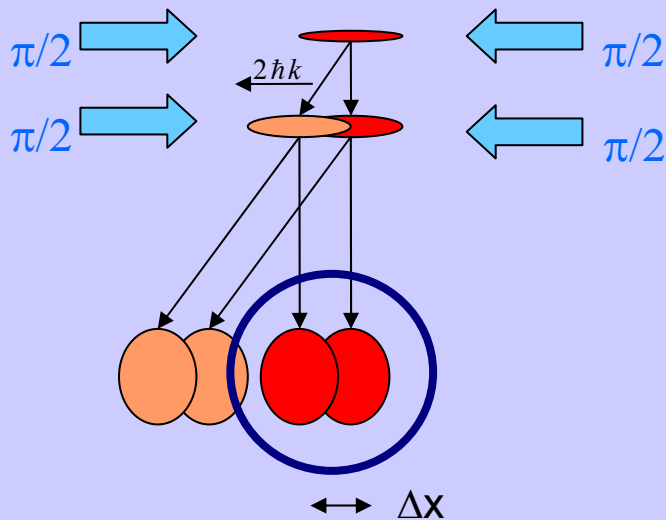
Measurement of intensity correlations
in one output port

= measurement of the spatial
second order correlation function
of the original condensate !

➡ coherence length

Advantage:
insensitive to global phase shifts

Interferometric scheme



Measurement of intensity correlations
in one output port

**= measurement of the second
order correlation function of the
original condensate !**

Normalised correlation function:

$$\gamma_f^{(2)}(x_1, x_2, \Delta x) = \frac{\langle (\hat{I}_1 - \langle \hat{I}_1 \rangle) (\hat{I}_2 - \langle \hat{I}_2 \rangle) \rangle}{\sqrt{\langle (\hat{I}_1 - \langle \hat{I}_1 \rangle)^2 \rangle \langle (\hat{I}_2 - \langle \hat{I}_2 \rangle)^2 \rangle}}$$

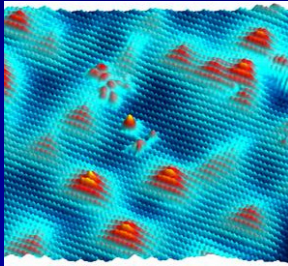
$$= \cos[\kappa(x_1 - x_2)] \exp\left[-\frac{\delta^2}{2} f^{(2)}(x_1, x_2, \Delta x)\right]$$

$$= \cos[\kappa(x_1 - x_2)] g^{(2)}(x_1, x_2, \Delta x)$$

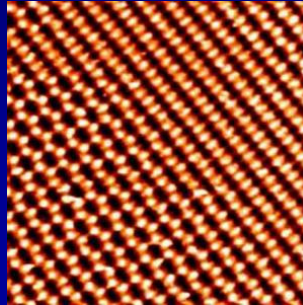
$$g^{(2)}(x_1, x_2, x_3, x_4) =$$

$$\exp\left[\frac{L}{2L_\phi} f^{(2)}(x_1, x_2, x_3, x_4)\right]$$

- disorder is present in various systems



www.omicron.de



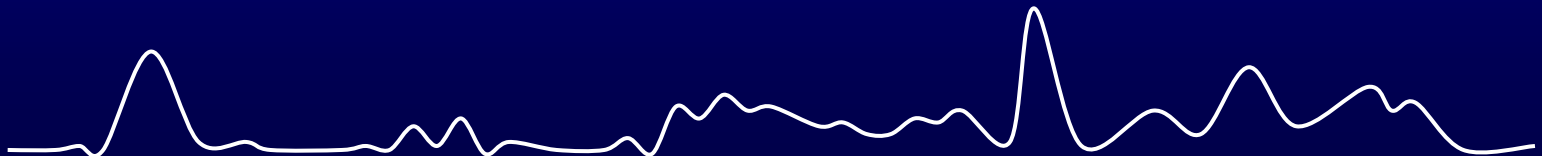
Suppression of superfluidity of ^4He in porous media with disorder.

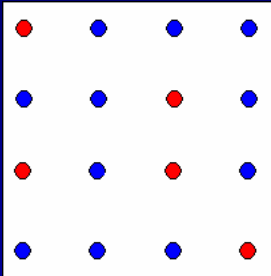
- drastic (non-perturbative) effects on physical properties (e.g. transport, optical)

$$H(\lambda) = H_0 + \lambda V(\vec{r}) \quad \text{non-perturbative}$$

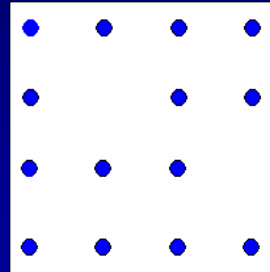
$$\forall \lambda : \lambda \neq 0 \quad \text{if } V(\vec{r}) \text{ is random / disordered}$$

⇒ Effect of controllable disorder on dynamics of quantum many-particle-systems?

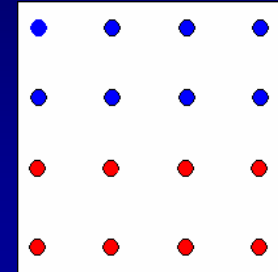




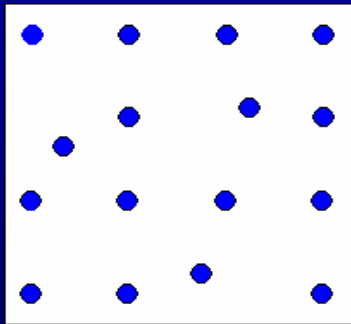
Mixtures



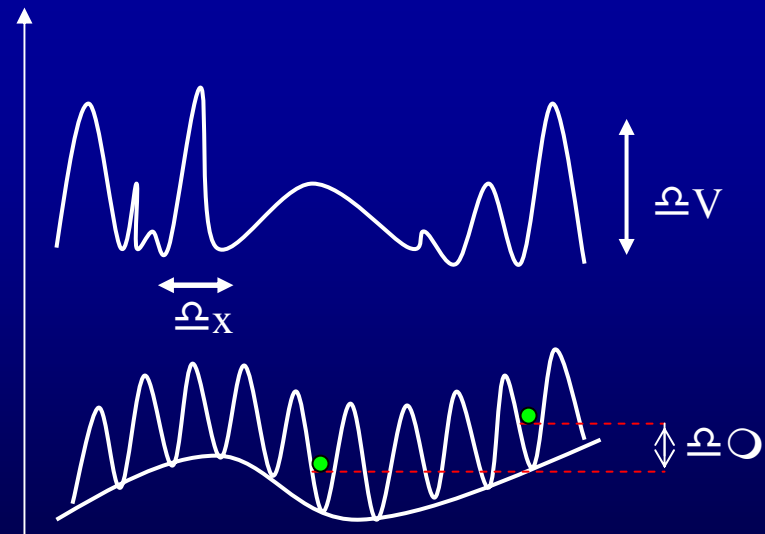
Empty sites



Surfaces, alloys



Lattice disorder

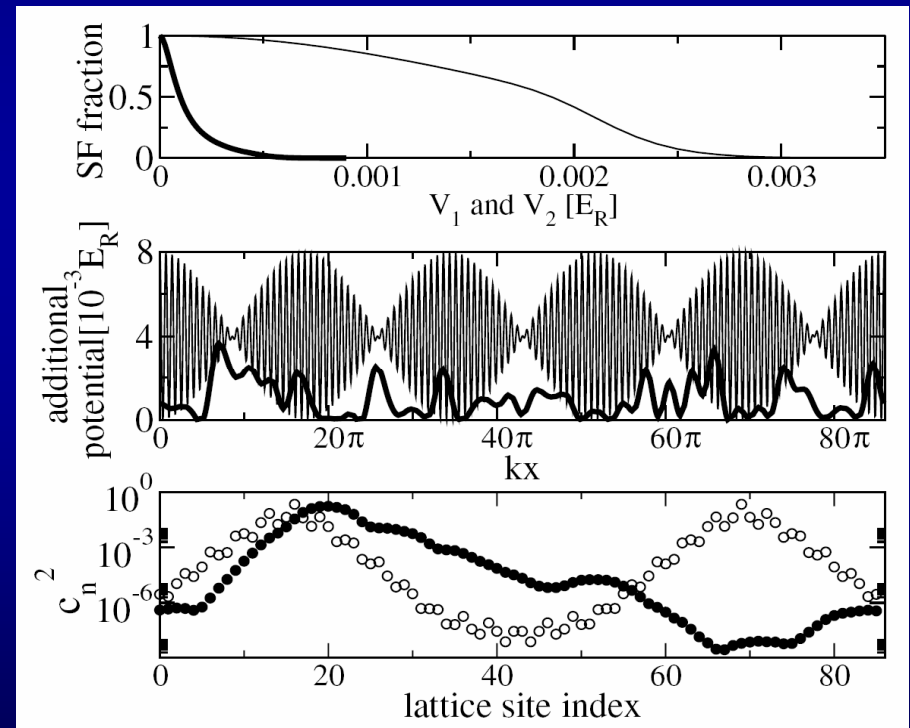
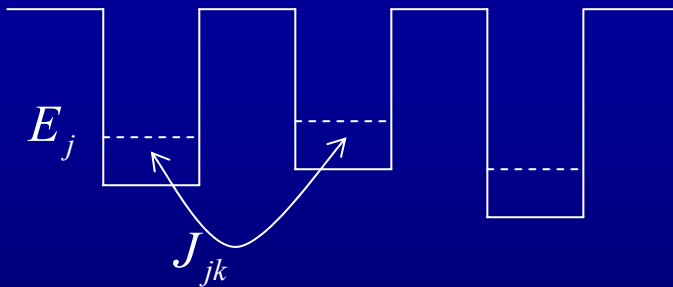


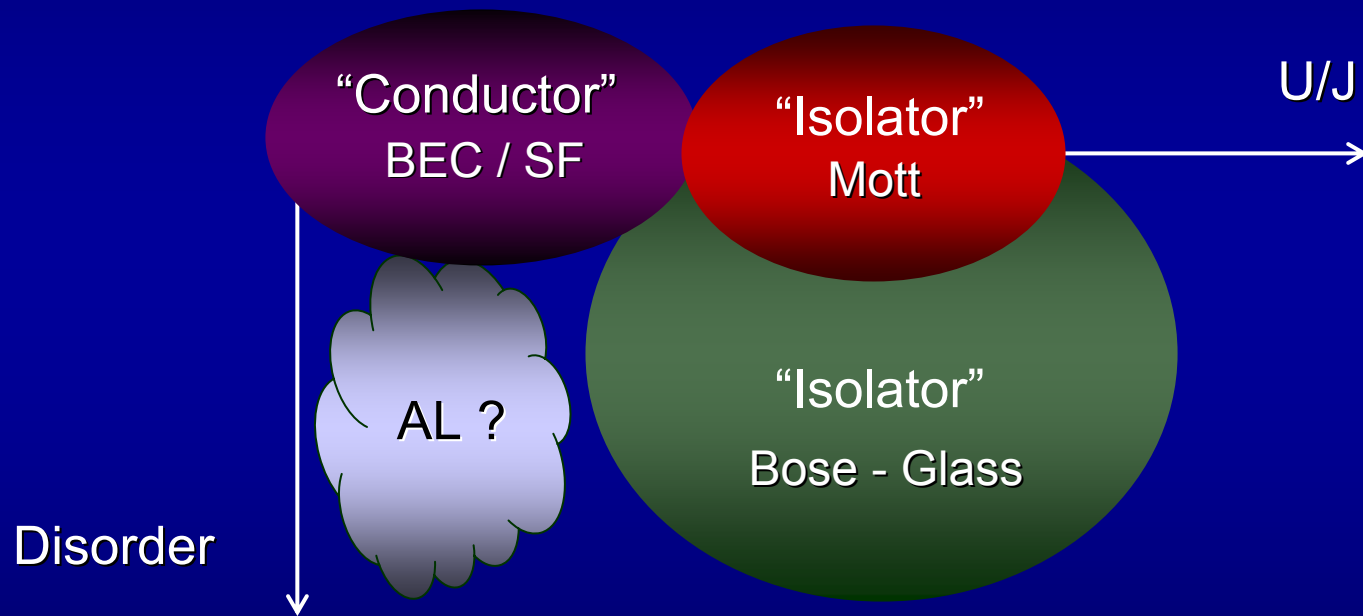
Non-interacting particles in 1D lattices :

- inhibition of transport , vanishing SF
- associated with localized states

Anderson model :

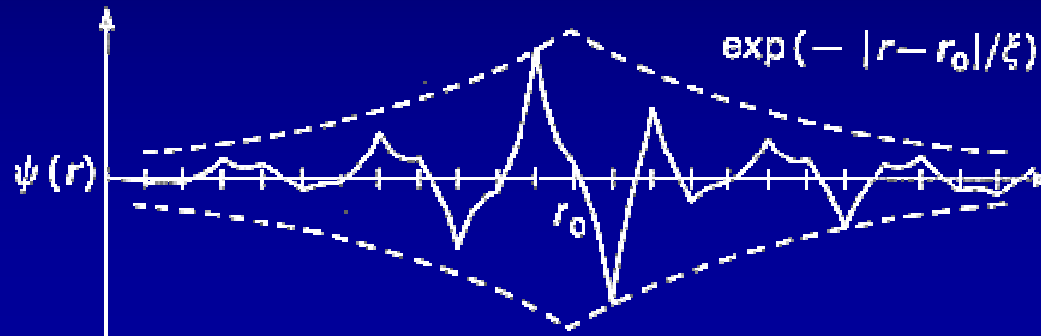
$$i \hbar \dot{a}_j = E_j a_j + \sum_{k \neq j} J_{jk} a_k$$





Can there be an Anderson Localization regime ?

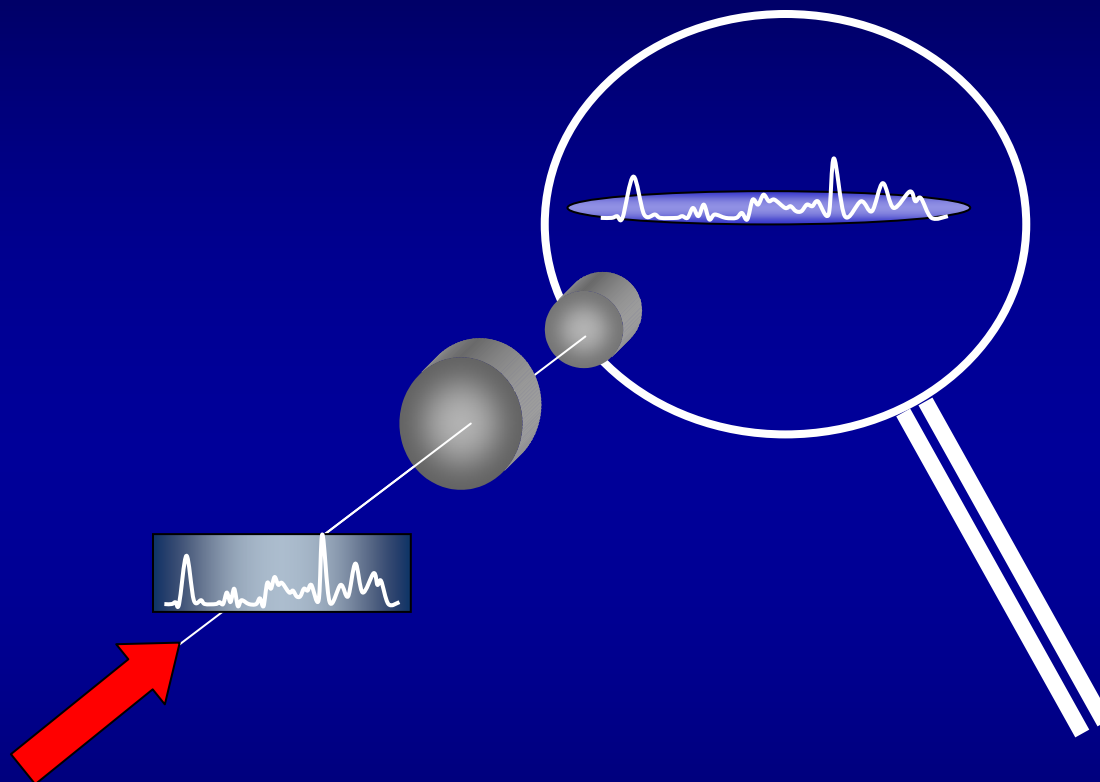
→ It strongly enhances the existence of localized states !



Wave function :
$$\psi(r) = \varphi(r) \exp\left(\frac{-|r-r_0|}{\xi}\right)$$
 ξ localization length

Anderson Localization is characterized by:

- vanishing superfluid fraction
- localization of atoms due to interference
- gapless excitation spectrum



this work:

T. Schulte et al.
Phys. Rev. Lett. 95, 170411 (2005),
cond-mat/0507453.

similar investigations:

J. E. Lye et al. Phys. Rev. Lett. 95, 070401 (2005)

C. Fort et al. Phys. Rev. Lett. 95, 170410 (2005)

D. Clément et al. Phys. Rev. Lett. 95, 170409 (2005)

Disorder potential :

Wavelength : 825 nm

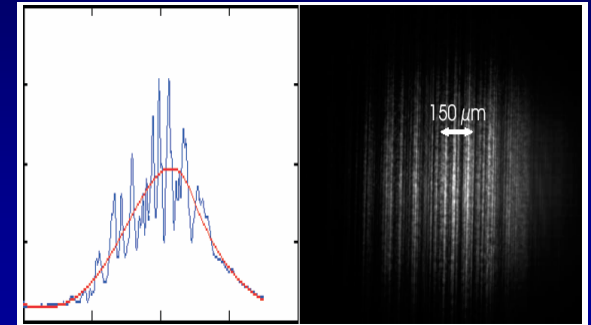
Waist : $\sim 480 \mu\text{m}$

Theo. modulation depth : $\sim 3 E_{\text{Rec}}$

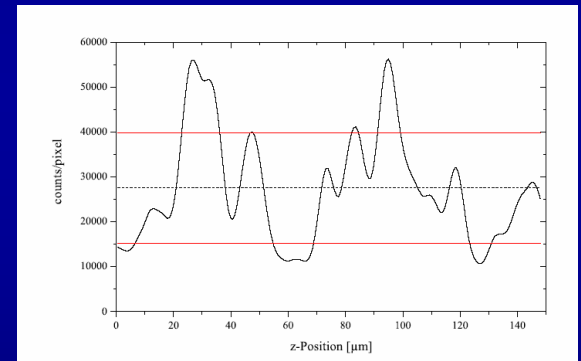
Smallest structure : $\sim 7 \mu\text{m}$

Modulations over cloud size : ~ 20

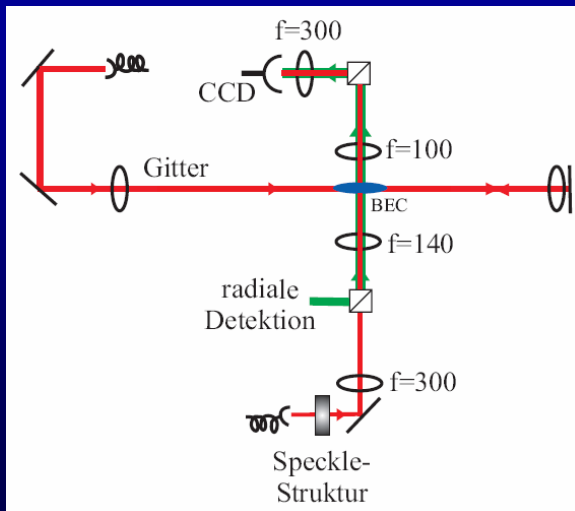
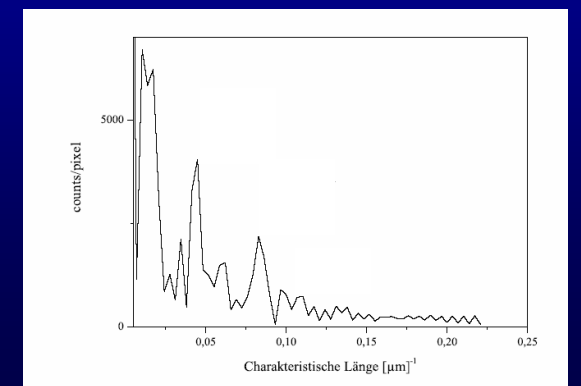
Image of total
beam profile



Intensity modulations
over cloud size



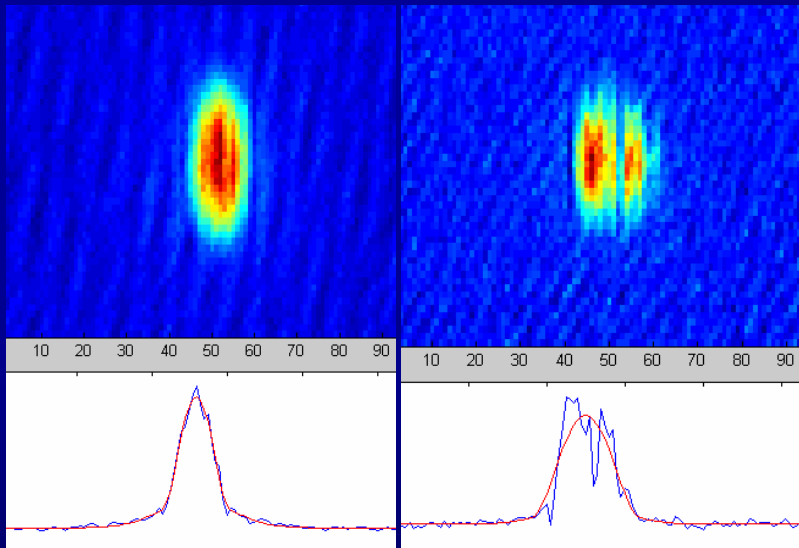
Fourier analysis
of profile



Absorption images :

TOF = 20.4 ms

Modulation depth $0.2 E_{rec}$



$N = 5 \cdot 10^4$

$N = 2.6 \cdot 10^4$

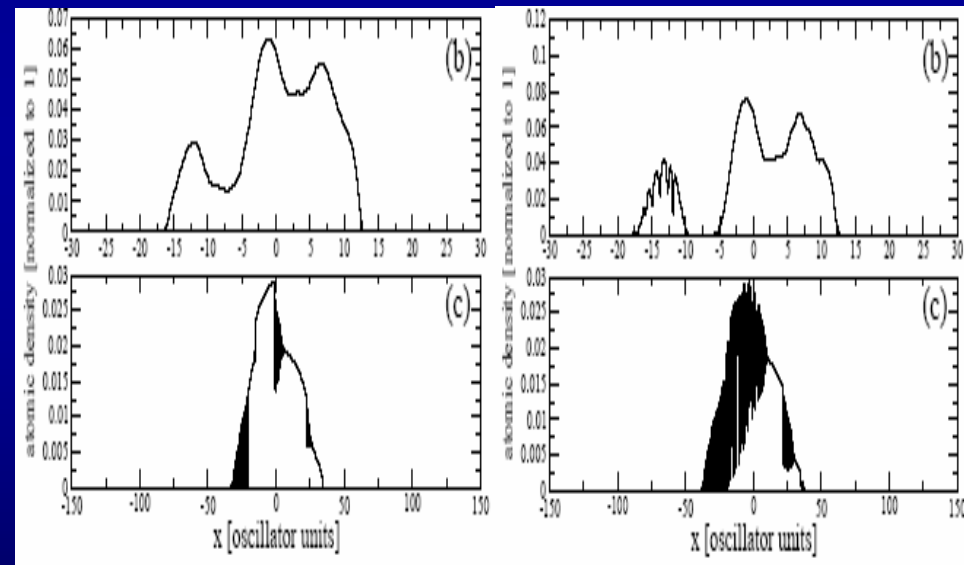
Numerical simulation :

TOF = 20.4 ms

Modulation depth

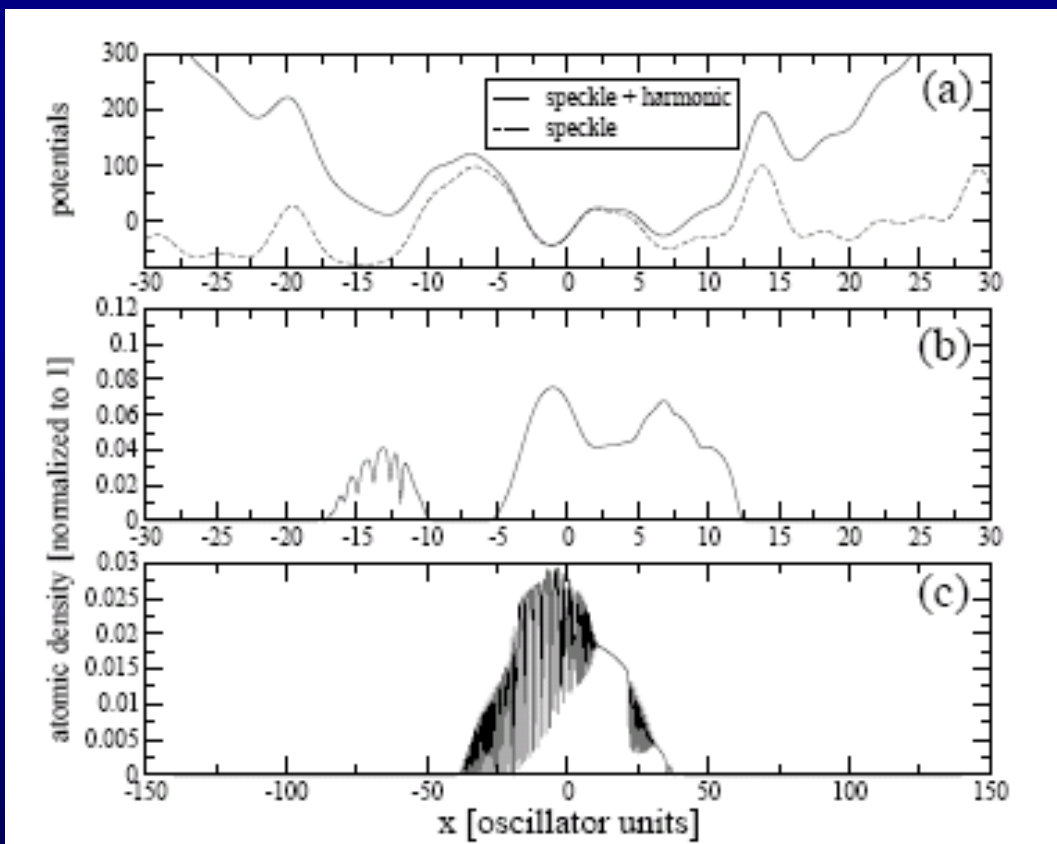
$0.2 E_{rec}$

$0.4 E_{rec}$



Parameters adjusted for proper TF - radius

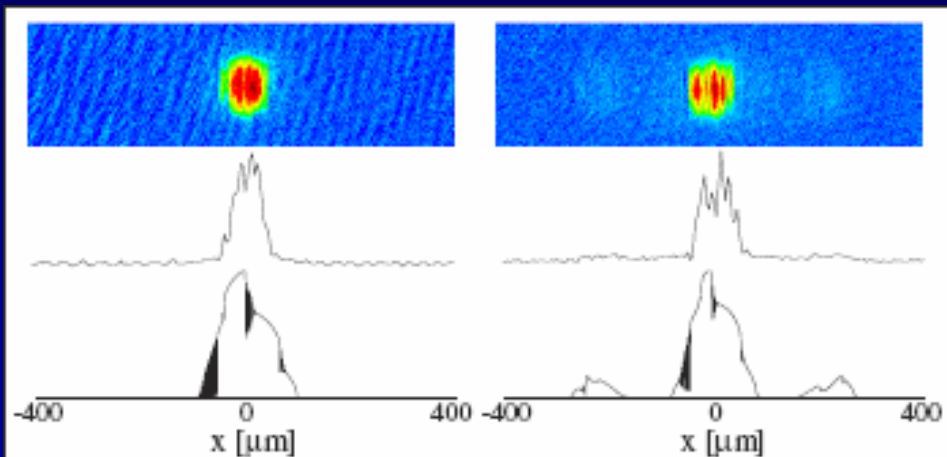
Numerical simulation : 1D GPE



Potential:
magnetic trap + disorder $0.4 E_r$

atomic density in trap

atomic density after 20.4 ms TOF



Experimental parameters:

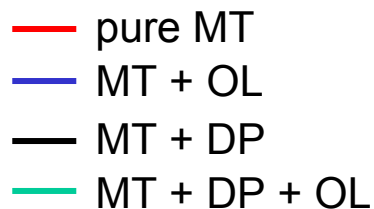
$$v_z = 14 \text{ Hz}; v_{\text{rad}} = 200 \text{ Hz}$$

$$N_{\text{BEC}} \sim 1,5 - 8 \cdot 10^4$$

axial width after 20.4 ms TOF :

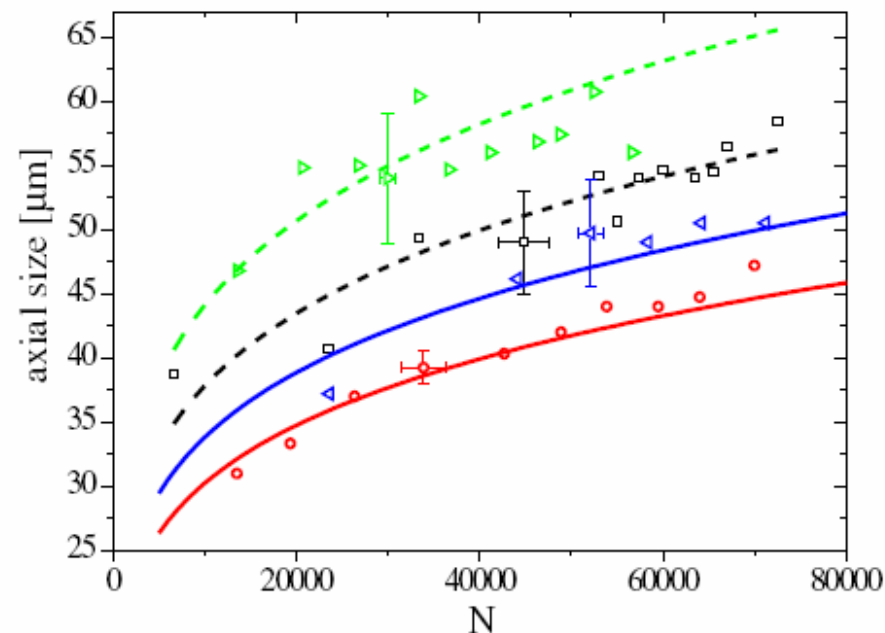
Modification of density profile:

- Pronounced fringes
- Axial expansion of ground state
- **no Anderson-localized regime**



Increase : 25 %

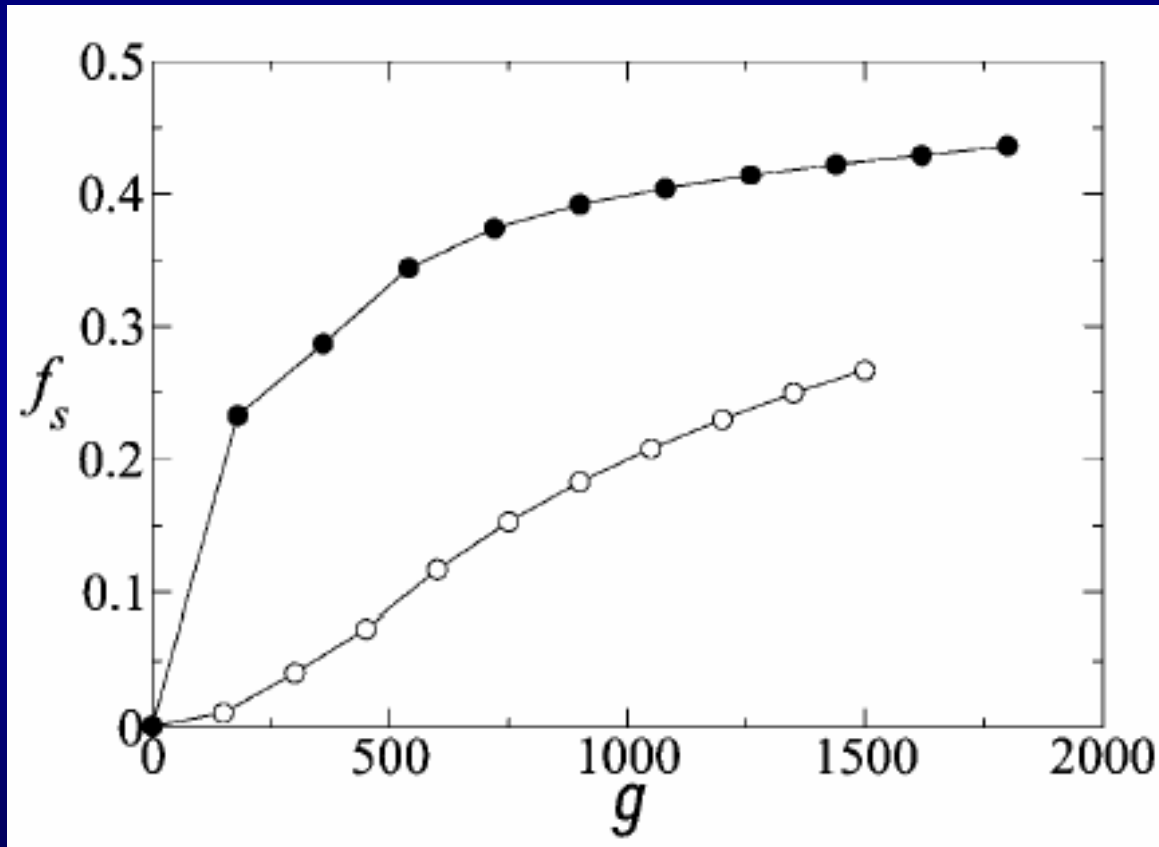
Increase : 28 %



Experimental difficulties in detecting AL-regime

LEB 175

Universität Hannover



full: $\nu_z = 14$ Hz

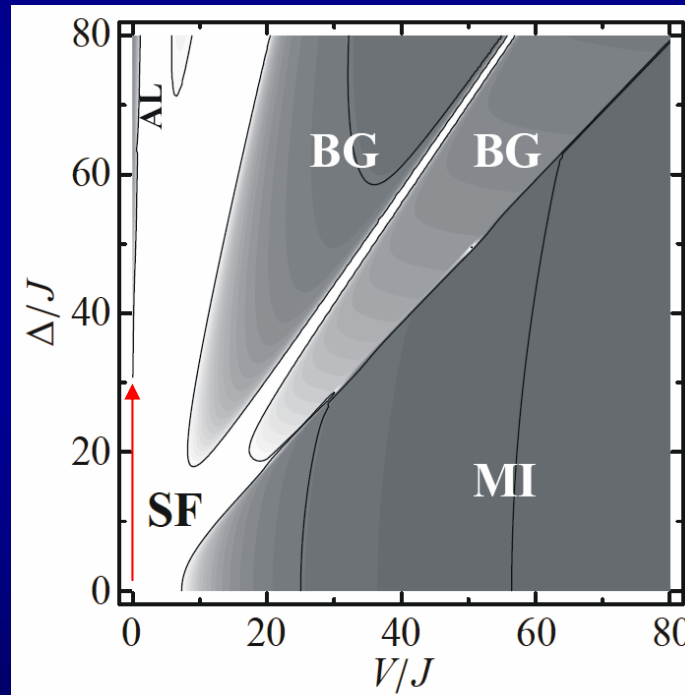
open: $\nu_z = 4$ Hz

ANDERSON LOCALISATION

1. Long-range phase coherence
2. High number fluctuations
3. gapless excitation spectrum

BOSE-GLASS PHASE

1. No phase coherence
2. Low number fluctuations
3. continuous excitation spectrum



SUPERFLUID PHASE

1. Long-range phase coherence
2. High number fluctuations
3. continuous excitation spectrum

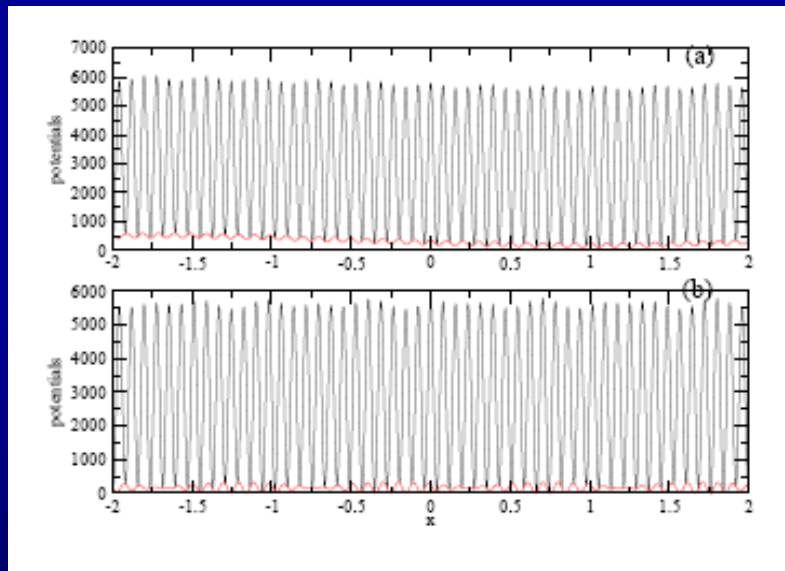
(R. Roth and K. Burnett,
PRA 68, 023604 (2003))

MOTT INSULATOR PHASE

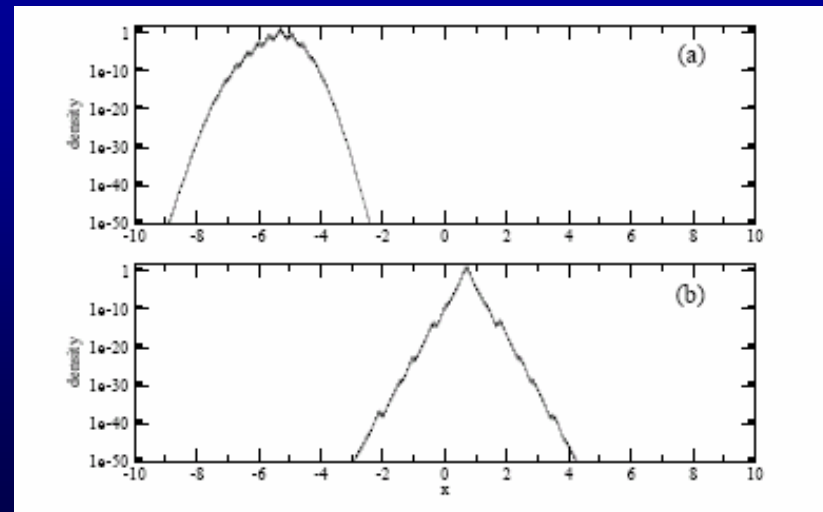
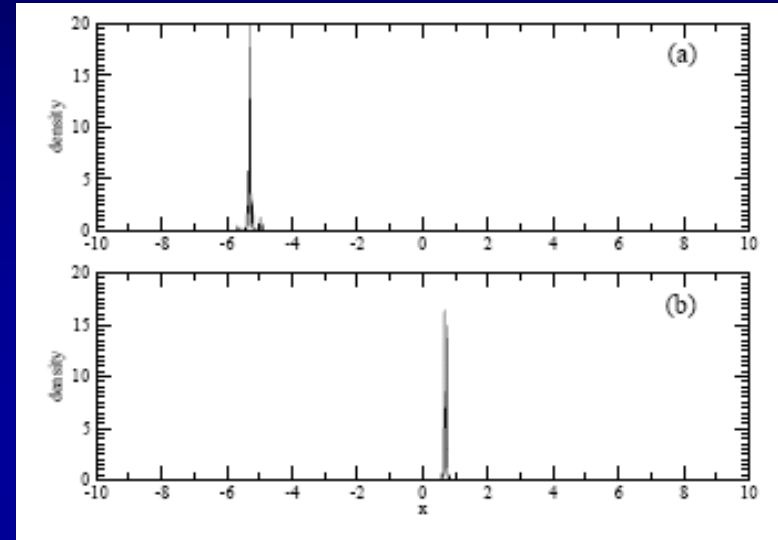
1. No phase coherence
2. Zero number fluctuations
3. discrete excitation spectrum

Simulation using **small scale disorder without interactions** to find localization regime.

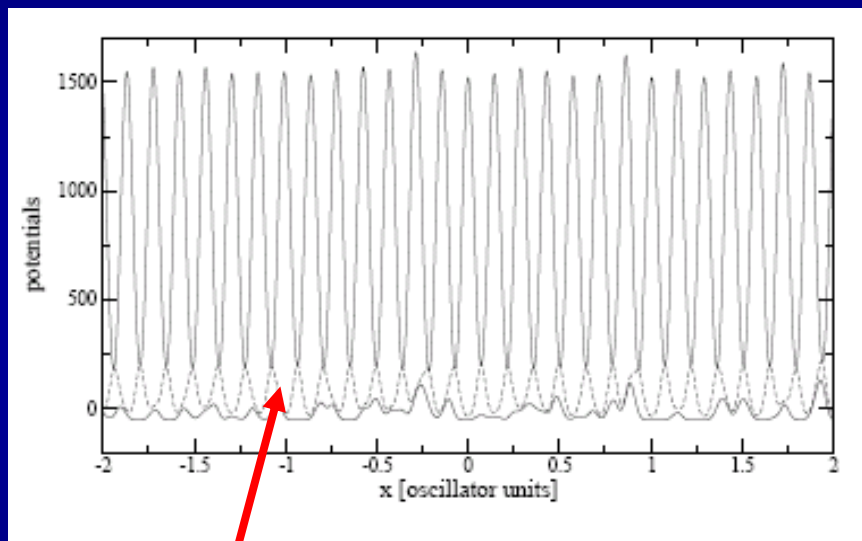
$$g=0$$



- a) Speckle pattern.
- b) Pseudorandom potential 1060nm + 960nm

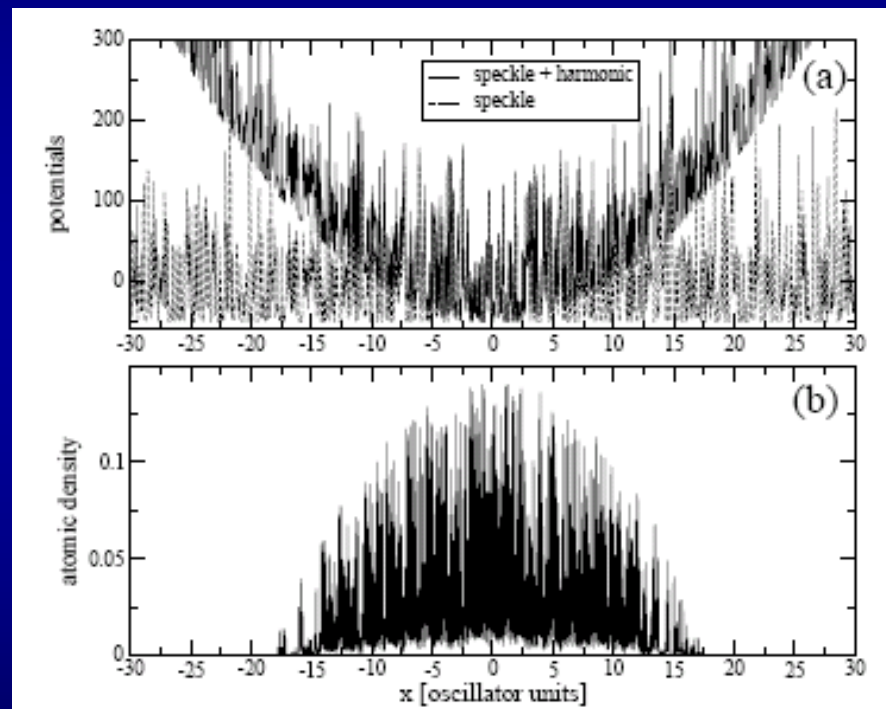


Simulation using **small scale disorder** to find localization regime



$$V_{\text{speckle}}(x) + g|\phi_0|^2$$

Smoothing of the potential
due to interactions!

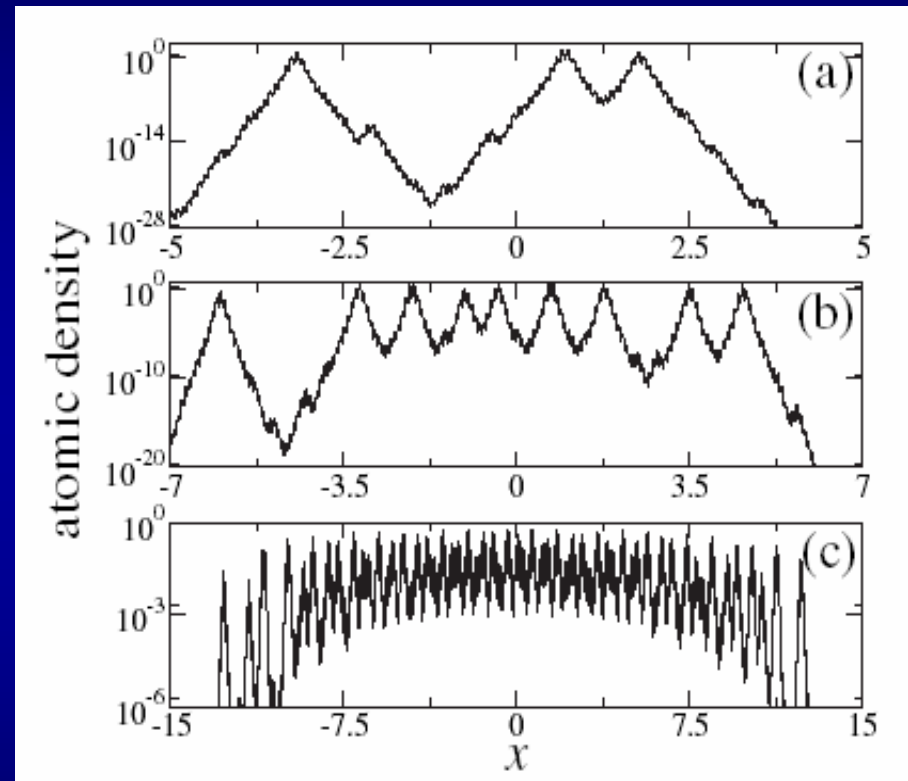


Is there localization?

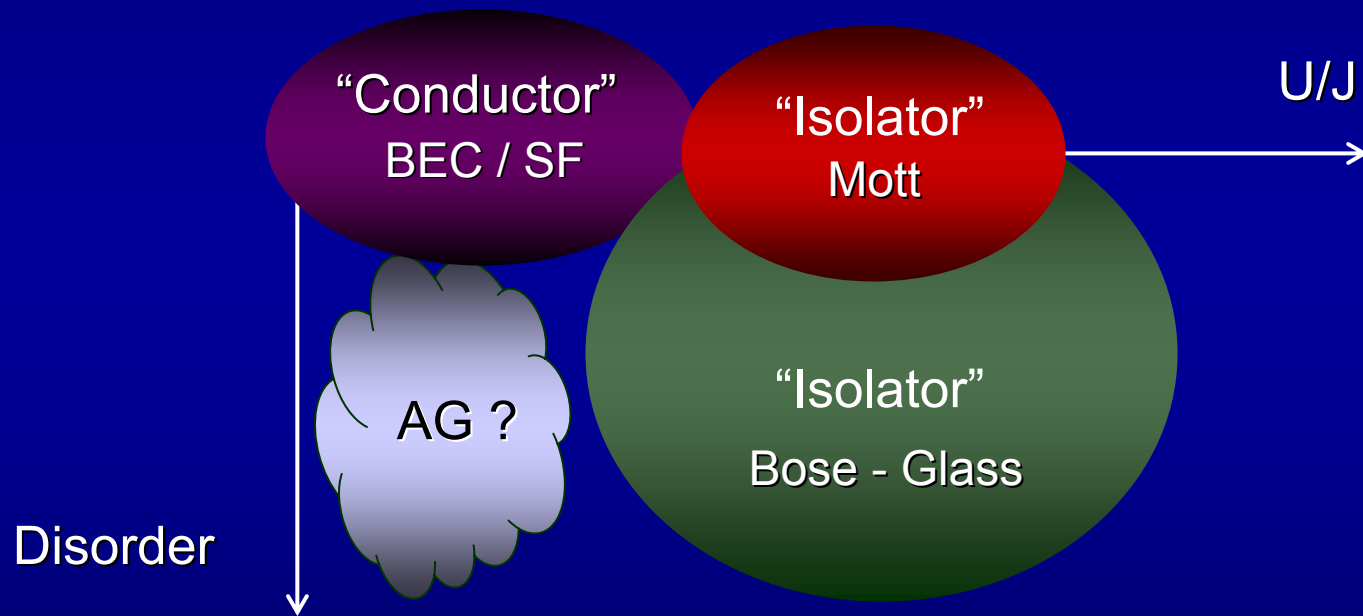
$g=8$

$g=128$

$g=256$



Are there experimentally reasonable parameters to observe the AL-regime?



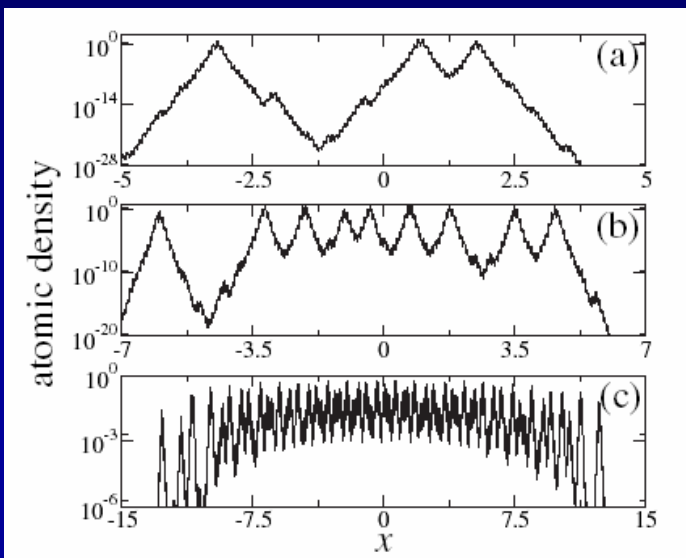
Can there be an Anderson Localization regime ?

Theoretical analysis of interaction-dependence

$g=8$

$g=128$

$g=256$



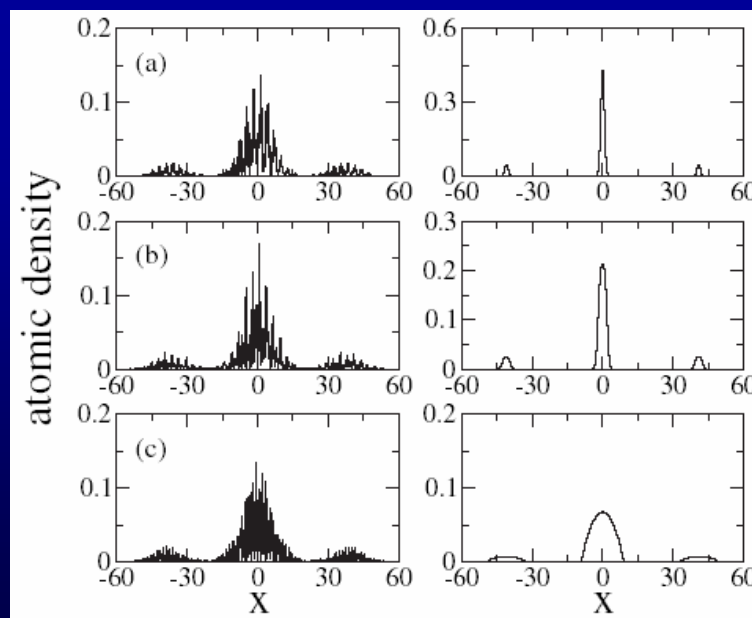
Are there experimentally reasonable parameters to observe the AL-regime?

$$v_z = 4 \text{ Hz}; \quad v_{\text{rad}} = 40 \text{ Hz}$$

$$g = 256$$

Free (but limited) parameters:

- Lowering interaction-strength g by Feshbach-resonance
- Variation of trap frequencies



$g=8$

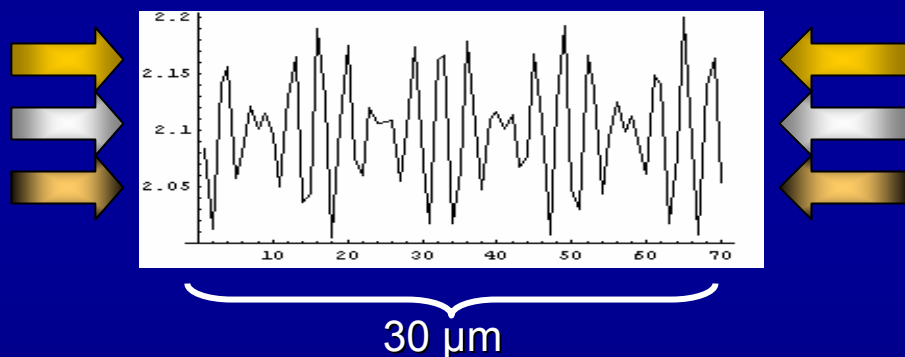
$g=128$

$g=256$

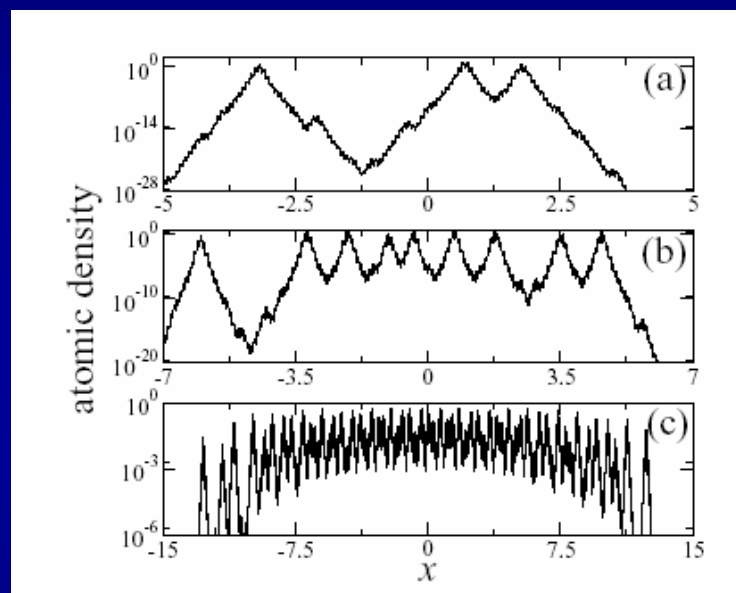
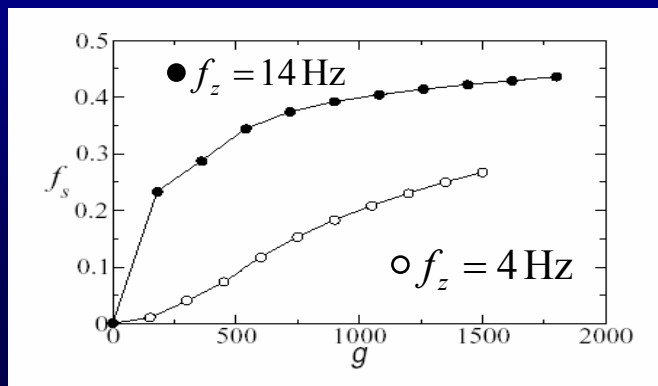
Towards Anderson localization :

• Realization of fine scaled disorder

- 2 incommensurate super lattices
e.g. @ 1040 nm + 980 nm



• Reduce interactions



a) $g = 0.5$

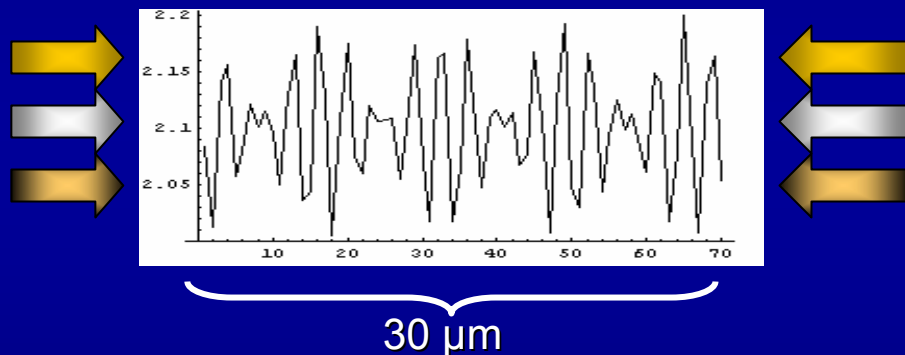
b) $g = 8$

c) $g = 256$

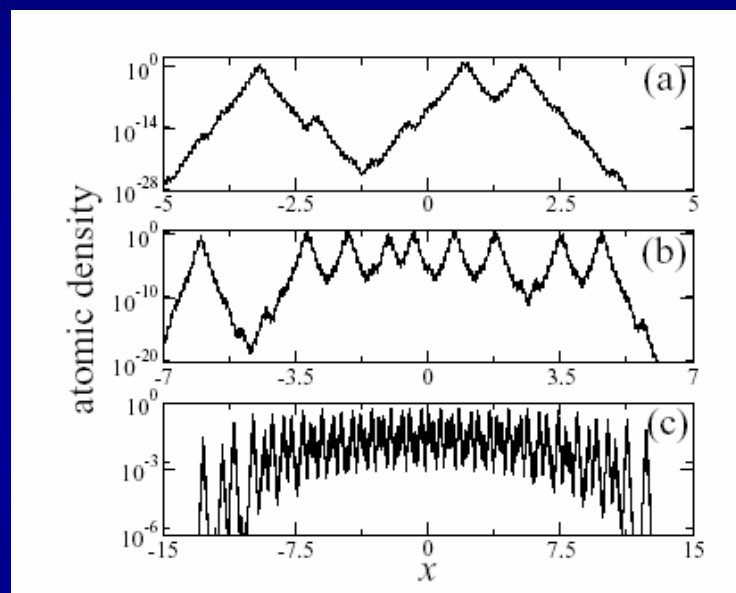
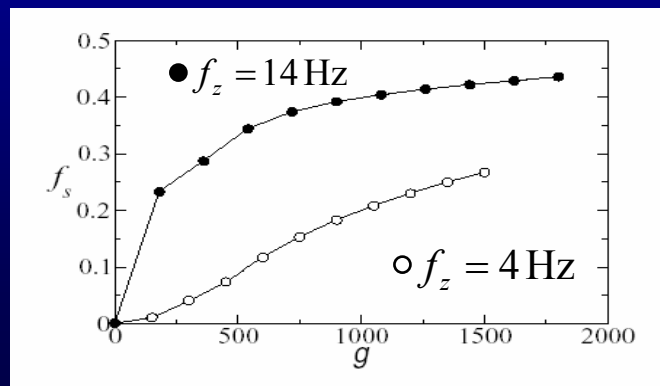
Towards Anderson localization :

• Realization of fine scaled disorder

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e.g. @ 1040 nm + 980 nm



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a) $g = 0.5$

b) $g = 8$

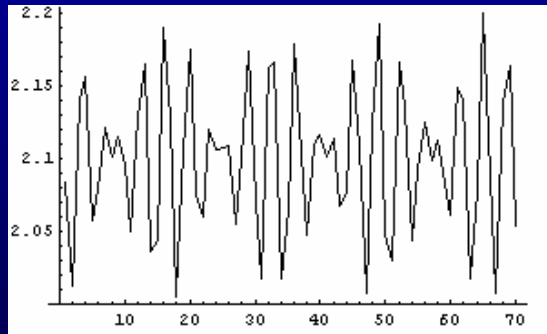
c) $g = 256$

→ Corresponding 3D trap :

$f_\rho = 40 \text{ Hz}$, $f_z = 4 \text{ Hz}$, $N = 10.000$

Anderson Localization :

- decrease mean density
 - work at small U/J
 - use appropriate disorder potential
- 2 incommensurate super lattices
e.g. @ 1040 nm + 980 nm

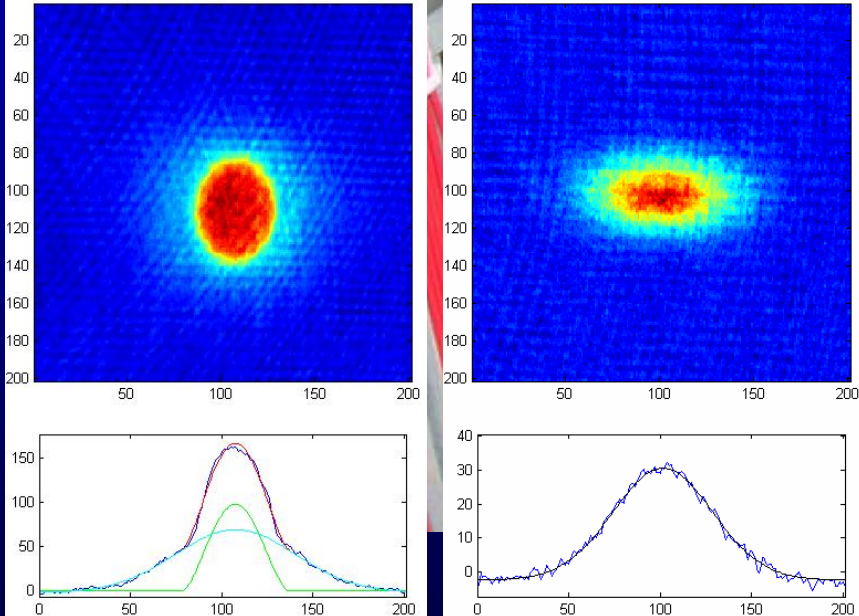
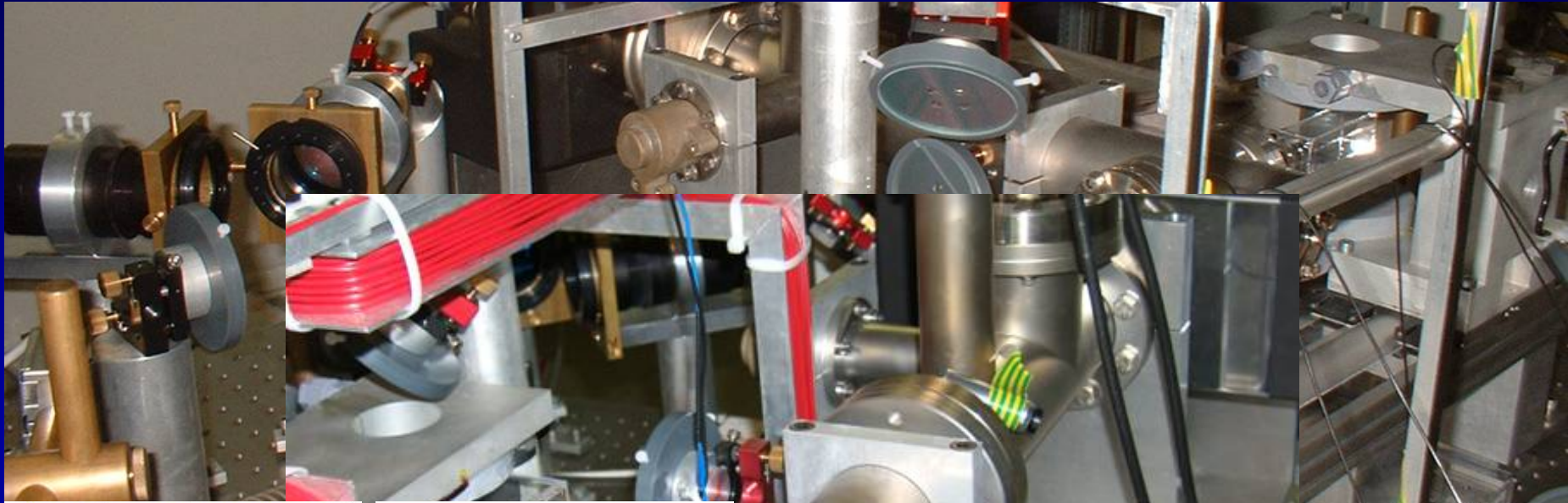


30 μm

...towards Bose Glass :

- high mean density
 - work at large U/J
 - find appropriate disorder potential
- 2 incommensurate super lattices
e.g. @ 1040 nm + 980 nm

A quantum degenerate Boson (^{87}Rb) and Fermion (^{40}K) - mixture



BEC (TOF 10 ms)

$\sim 7 \times 10^5$ atoms

$T < 200$ nK

$T/T_C < 0.3$

QDF (TOF 2 ms)

$\sim 1.5 \times 10^5$ atoms

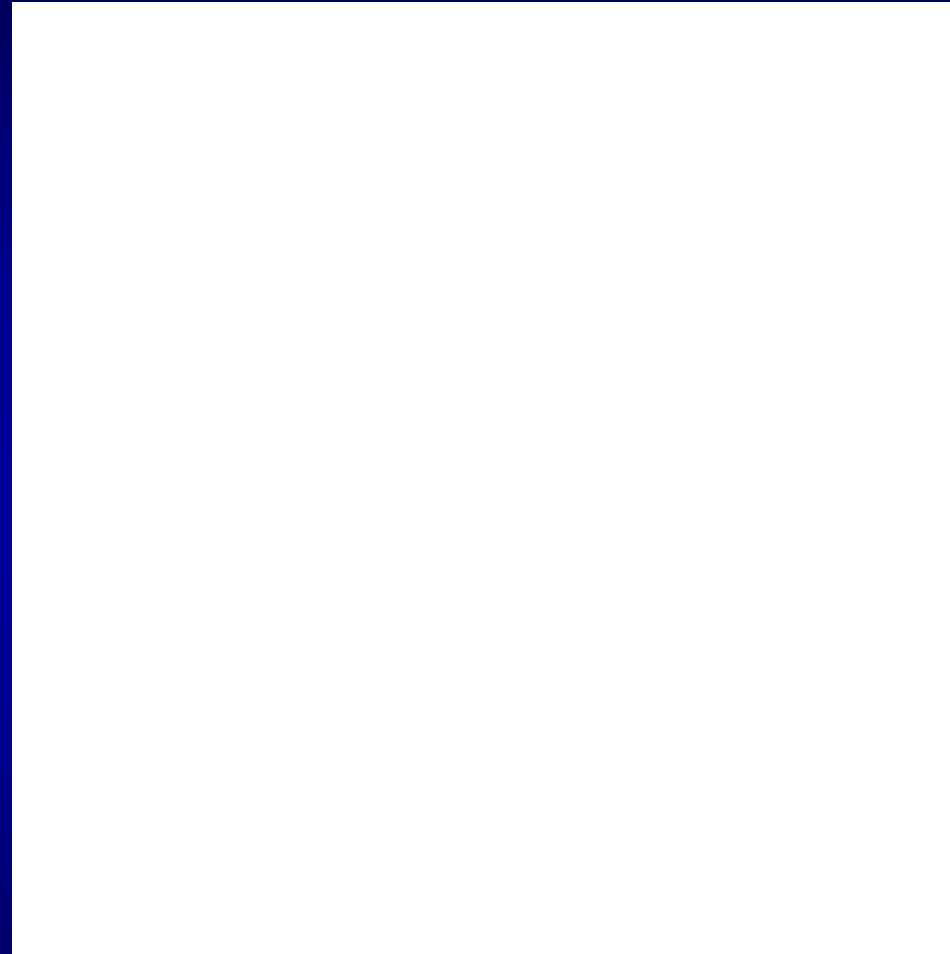
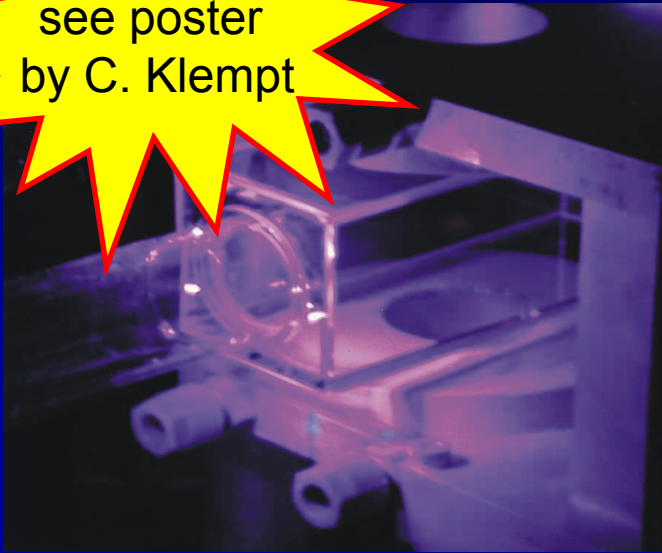
$T < 200$ nK

$T/T_F < 0.3$

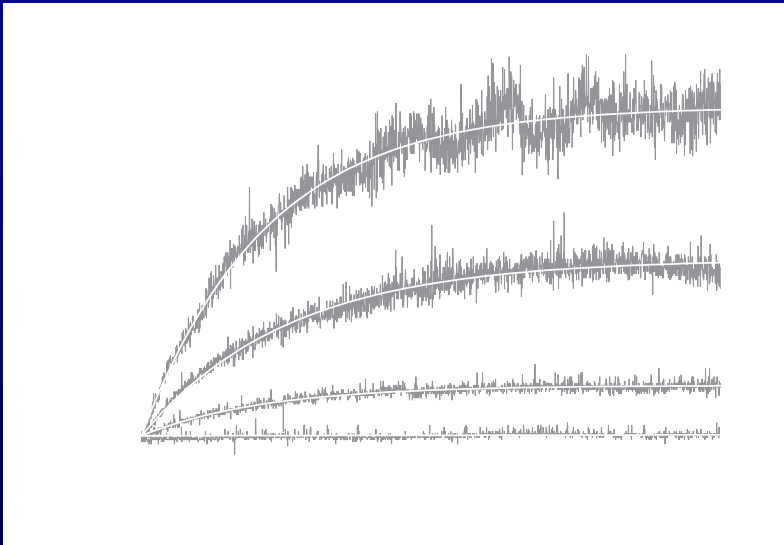
new in
Hannover:

UV light-induced atom desorption (LIAD) for large rubidium and potassium MOTs

see poster
by C. Klempt



LIAD intensity dependence



Number of atoms in a Rb MOT for various LIAD wavelengths

C. Klempt, T. van Zoest, T. Henninger, O. Topic, E. Rasel,
W. Ertmer, and J. Arlt, Phys. Rev. A 73, 013410 (2006)