

Atom-Light Entanglement

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Experiments :

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Objectives

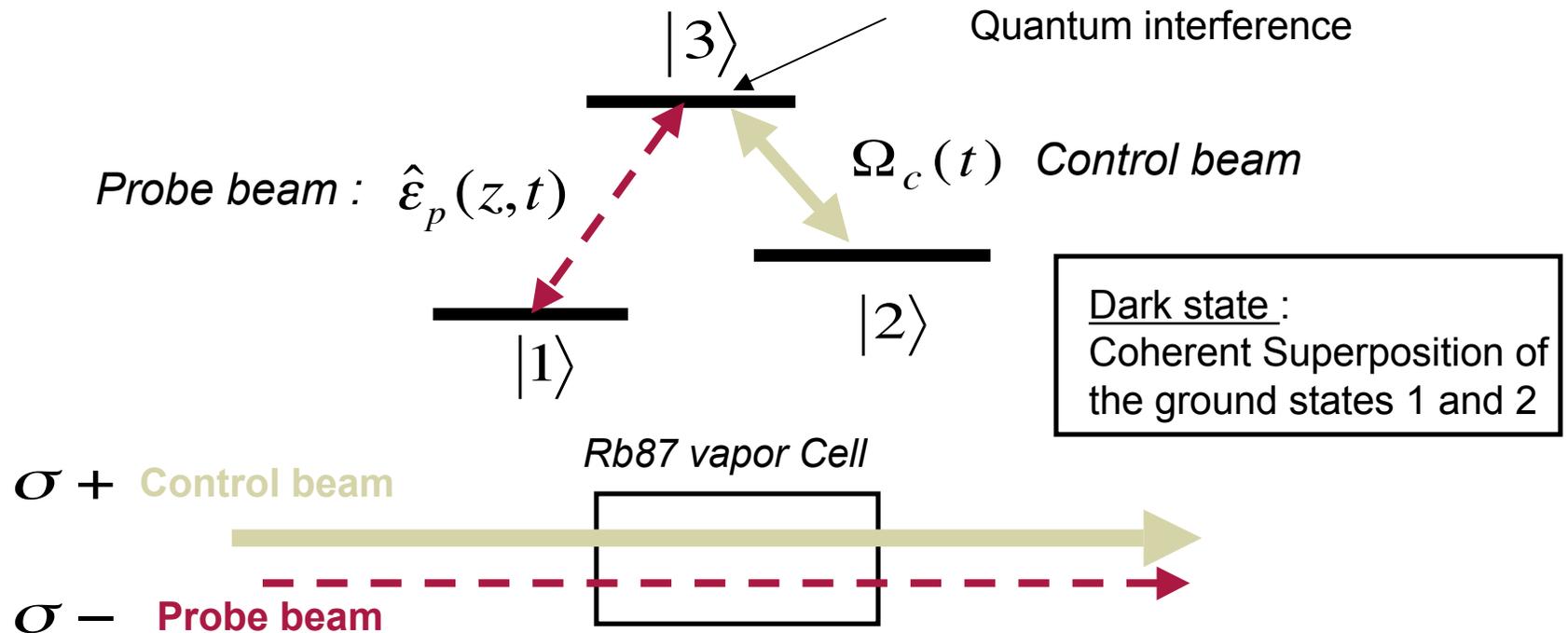
→ **Quantum information storage**

AM/PM (Conj. Obs.)
Squeezing
Entanglement

- Quantify EIT-based Quantum memories
 - Theoretically → Stochastic Simulations of light storage
 - Experimentally → Quantum Information Delay/Storage in atomic ensemble.
- Generate non-classical light at atomic wavelengths
 - Self rotation → Too noisy (*M.Hsu et al. To be published in PRA*)
 - SHG/OPO @ 795 nm → In progress

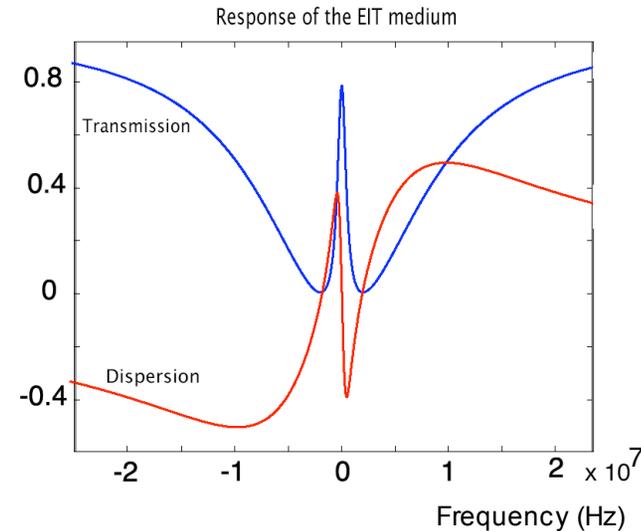
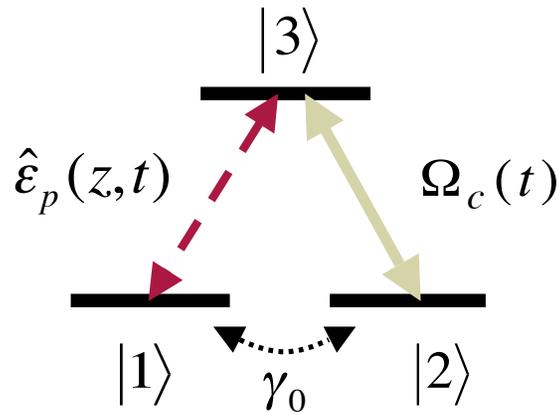
Transfer of quantum states between optical and atomic fields

Tool : Electromagnetically induced transparency (EIT)



We want to consider the interaction of the two **quadratures** of the probe field when interacting with such a system in **single pass** and measure their **variances**.

Quantum information delay



Parameters :

$$\left\{ \begin{array}{l} \Omega_c(t) = 5\gamma \\ \gamma_0 = 1\text{kHz} \\ N = 1.10^{12} \text{at.cm}^{-3} \\ \Delta\nu = 2\text{MHz} \\ \tau_d = 10\mu\text{s} \end{array} \right.$$

Langevin treatment of the Lambda system in the weak probe approx. :

$$\underline{V_{out}^\theta(\omega)} = \eta(\omega) \underline{V_{in}^\theta(\omega)} + (1 - \eta(\omega)) \underline{V_\nu(\omega)}$$

Input/Output quadrature variances

Vacuum field

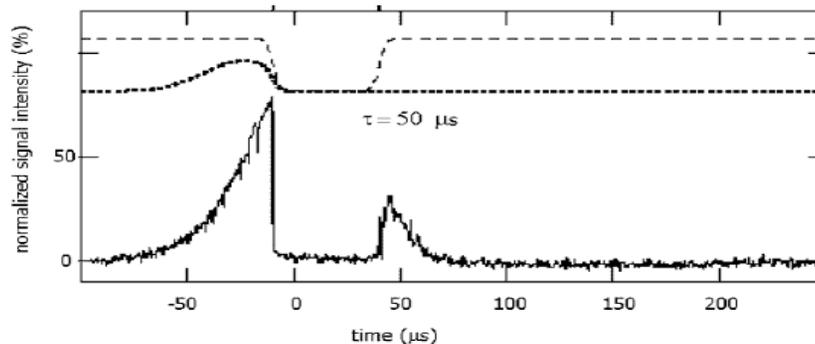
Beam splitter like relation : No extra atomic noise in this regime.

For example : With 4 dB of squeezing input to the system : $V^{sqz}_{out}(\omega) = 3.8\text{dB}$

Quantum Information Storage

- Switching off the coupling field : the probe field information is stored into the long lived coherence between the ground states when being delayed
- Switching on the coupling field : the probe field is retrieved

In vapor cells...

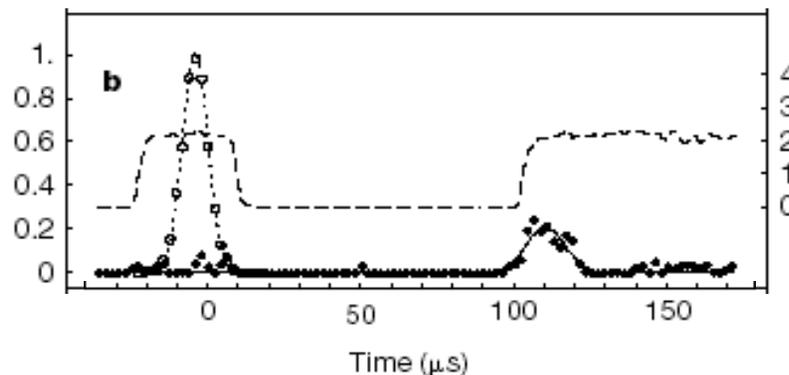


Phillips et al.
Phys. Rev. Lett., 86:783,
(2001).

All the experiments in vapor cells present so far the same exponentially decaying output pulse....

Controversy on the interpretation.

in MOTs...



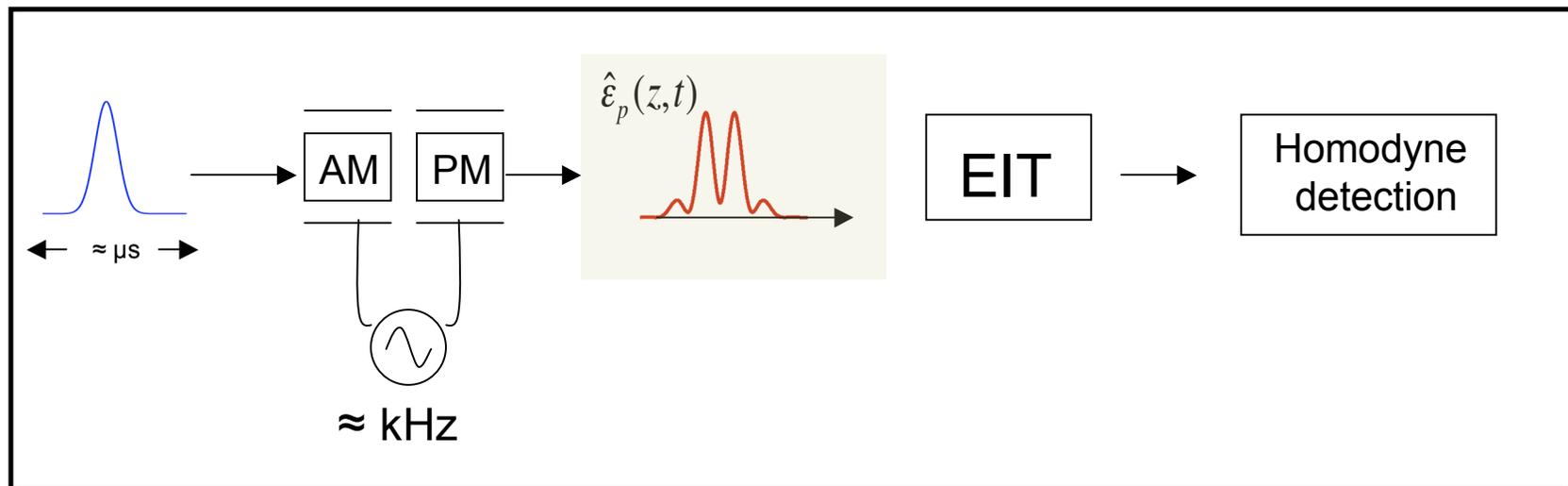
Liu et al.
Nature, 409:490, (2001).

Smaller decoherence rate.
Longer storage time

Quantum Information Storage

- Could the pulse shape be preserved ?
- What is the conjugate observable of a pulse shape ?
- How much quantum information can be stored ?
- Is it a noisy process ?

What we want to do :



Quantum Information Storage

Quantification of this process using quantum information tools:

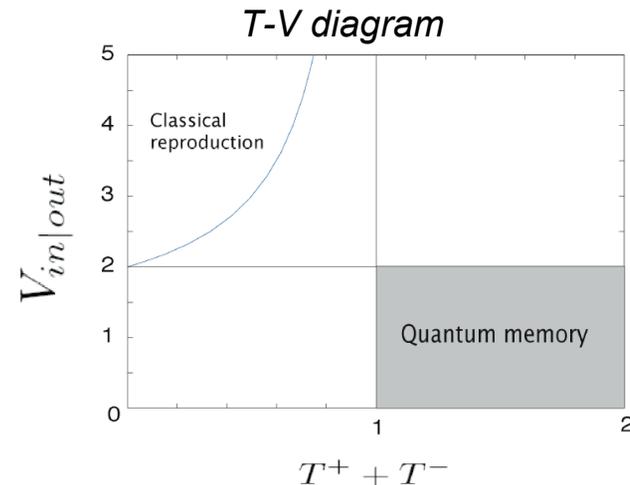
Similar to Q.teleportation ←

- Average fidelity : $\mathcal{F} = \langle \Psi_{in} | \rho | \Psi_{in} \rangle$

Bounded by 1/2 for a classical measure and prepare protocol.
Hammerer et al. (PRL 2005)

- Signal to noise transfer and conditional variances (T and V).
Useful when dealing with non-unity gain

Ralph and Lam criteria (PRL 1998)



Is it a Quantum Memory when taking into account realistic experimental parameters ?

Modeling using Stochastic methods

Interaction Hamiltonian for an optically thick medium :

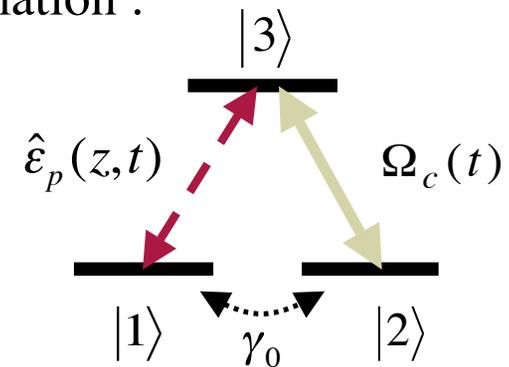
$$\hat{\mathcal{H}}_{\text{int}} = i\hbar n A_{\text{eff}} \int_0^l dz \left[-g \left(\hat{\mathcal{E}}_p(z, t) \hat{\sigma}_{13}(z, t) + \hat{\mathcal{E}}_p^\dagger(z, t) \hat{\sigma}_{31}(z, t) \right) + \Omega_c(t) \hat{\sigma}_{32}(z, t) + \Omega_c^*(t) \hat{\sigma}_{23}(z, t) \right]$$

Maxwell equations in the rotating wave approximation :

$$\begin{cases} \left(\frac{\partial}{\partial t} + c \frac{\partial}{\partial z} \right) \bar{\alpha}(z, t) = -gN \sigma_{13}^-(z, t) \\ \left(\frac{\partial}{\partial t} + c \frac{\partial}{\partial z} \right) \bar{\beta}(z, t) = -gN \sigma_{31}^-(z, t) \end{cases}$$

Master equation :

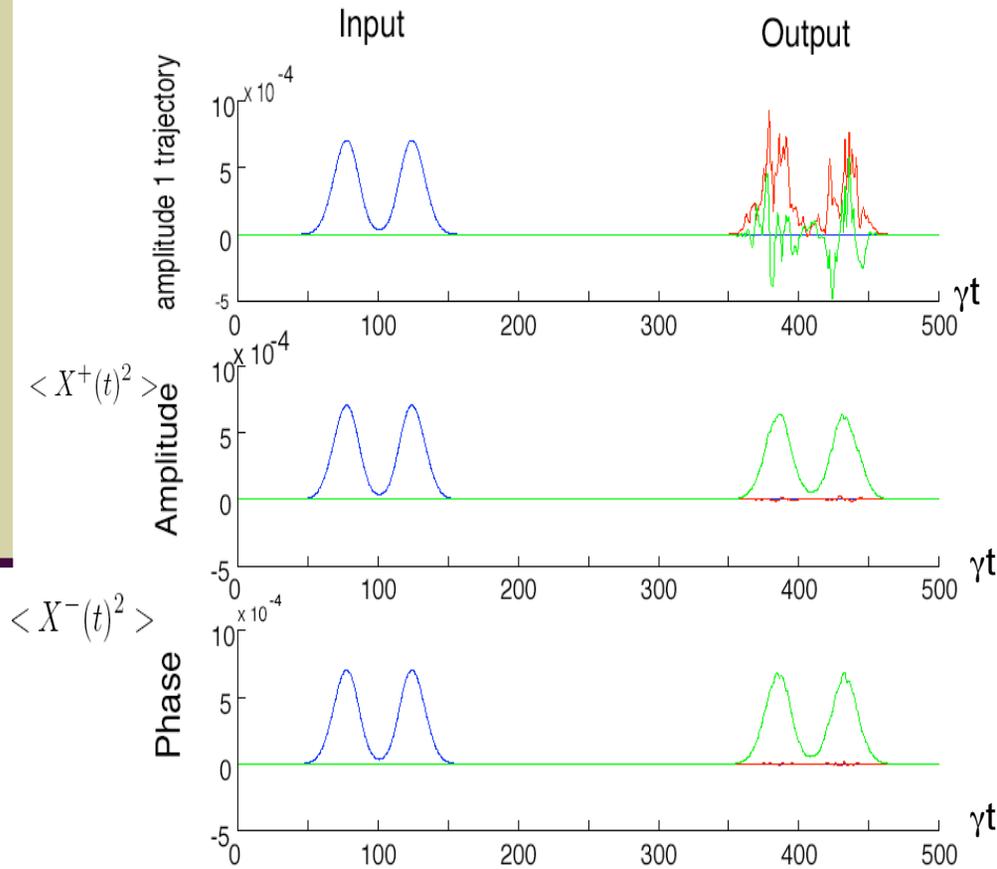
$$\frac{\partial}{\partial t} \hat{\rho} = i \hbar [\hat{\mathcal{H}}_{\text{int}}, \hat{\rho}] + \mathcal{L}_{13}[\hat{\rho}] + \mathcal{L}_{23}[\hat{\rho}] + \mathcal{L}_{[1,2]}[\hat{\rho}]$$



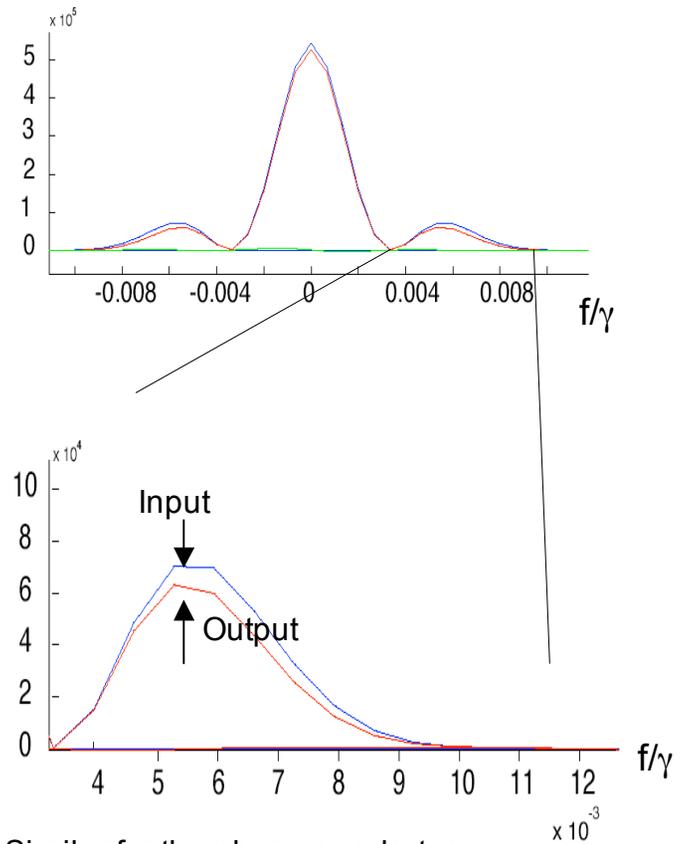
We use the *Positive P* representation,
define normal ordering and obtain a *Fokker Planck equation*.
We finally convert the FPE into a set of *SDE* (18 noise terms...)

Numerical results

Time domain



Fourier domain : $\langle X^+(\omega)X^+(-\omega) \rangle$



Similar for the phase quadrature...

Realistically...

Next step: Characterize the Quantum state transfer including decoherence factors:

$$\begin{aligned}\mathcal{L}_{[1,2]}[\hat{\rho}] = & \gamma_{coll} \sum_{z_k \in \delta z} (\hat{\sigma}_{12}^k \hat{\rho} \hat{\sigma}_{21}^k - \frac{1}{2} \hat{\sigma}_{12}^k \hat{\sigma}_{21}^k \hat{\rho} - \frac{1}{2} \hat{\rho} \hat{\sigma}_{21}^k \hat{\sigma}_{12}^k) \\ & + \gamma_{coll} \sum_{z_k \in \delta z} (\hat{\sigma}_{21}^k \hat{\rho} \hat{\sigma}_{12}^k - \frac{1}{2} \hat{\sigma}_{21}^k \hat{\sigma}_{12}^k \hat{\rho} - \frac{1}{2} \hat{\rho} \hat{\sigma}_{12}^k \hat{\sigma}_{21}^k)\end{aligned}\quad \text{Collisions}$$

$$\begin{aligned}\mathcal{L}_{[1,2]}[\hat{\rho}] = & \gamma_{deph} \sum_{z_k \in \delta z} \left((\hat{\sigma}_{11}^k - \hat{\sigma}_{22}^k) \hat{\rho} (\hat{\sigma}_{11}^k - \hat{\sigma}_{22}^k) - \frac{1}{2} (\hat{\sigma}_{11}^k + \hat{\sigma}_{22}^k) \hat{\rho} - \frac{1}{2} \hat{\rho} (\hat{\sigma}_{11}^k + \hat{\sigma}_{22}^k) \right) \\ & + \gamma_{deph} \sum_{z_k \in \delta z} \left((\hat{\sigma}_{22}^k - \hat{\sigma}_{11}^k) \hat{\rho} (\hat{\sigma}_{22}^k - \hat{\sigma}_{11}^k) - \frac{1}{2} (\hat{\sigma}_{22}^k + \hat{\sigma}_{11}^k) \hat{\rho} - \frac{1}{2} \hat{\rho} (\hat{\sigma}_{22}^k + \hat{\sigma}_{11}^k) \right)\end{aligned}\quad \text{Dephasing}$$

Main theoretical issue: Need for computing time to calculate the noise.

⇒ requires a lot more trajectories ($\approx 10^6$)

Main experimental issue:

Need for a good delay (to fit the whole pulse inside the cell)

⇒ narrow transparency window.

⇒ Low frequency information (Squeezing, AM/PM).